

Polarized radiation transfer in multidimensional models of the solar atmosphere

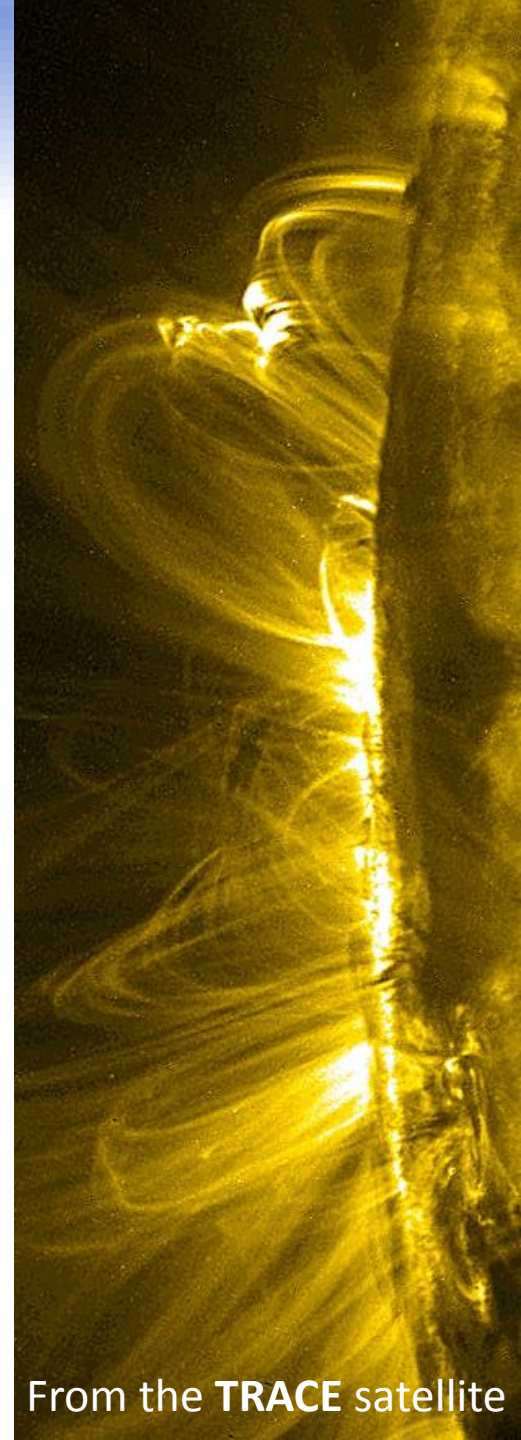
Tanausú del Pino Alemán

19th September 2019



Introduction

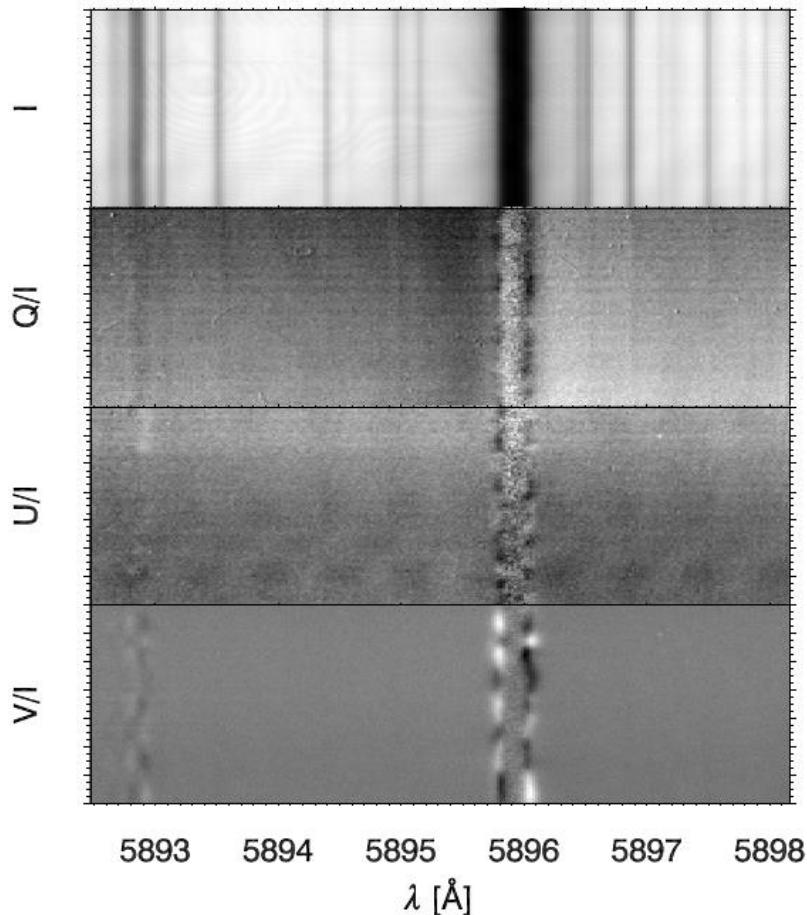
- **Magnetic fields** play a key role in the physics of the solar atmosphere
- Responsible of the **solar activity**
- Forms the **plasma structures** of the outer solar atmosphere
- Key to explain the existence of a hot (1 million degrees) **corona**



From the **TRACE** satellite

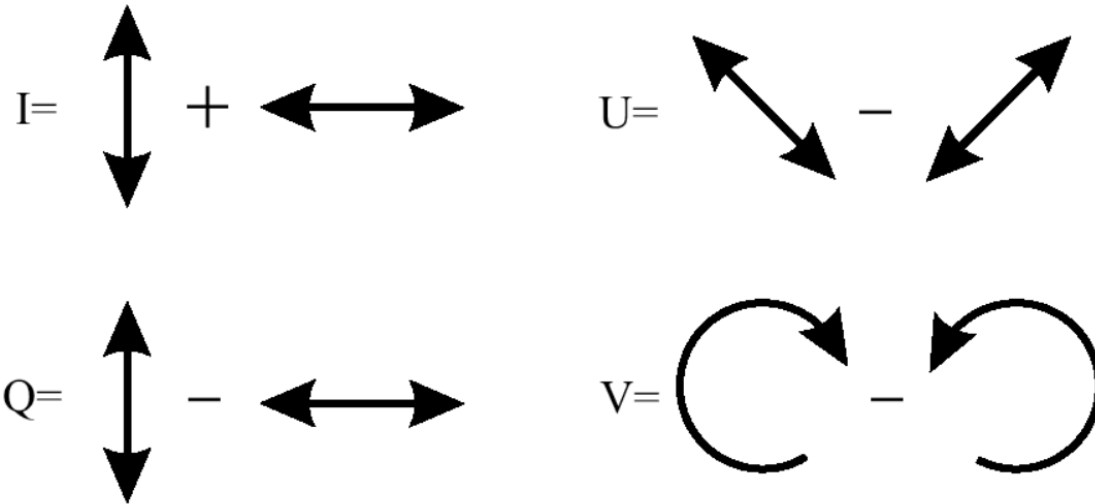
Introduction

Na I D1 slit observation with ZIMPOL@IRSOL



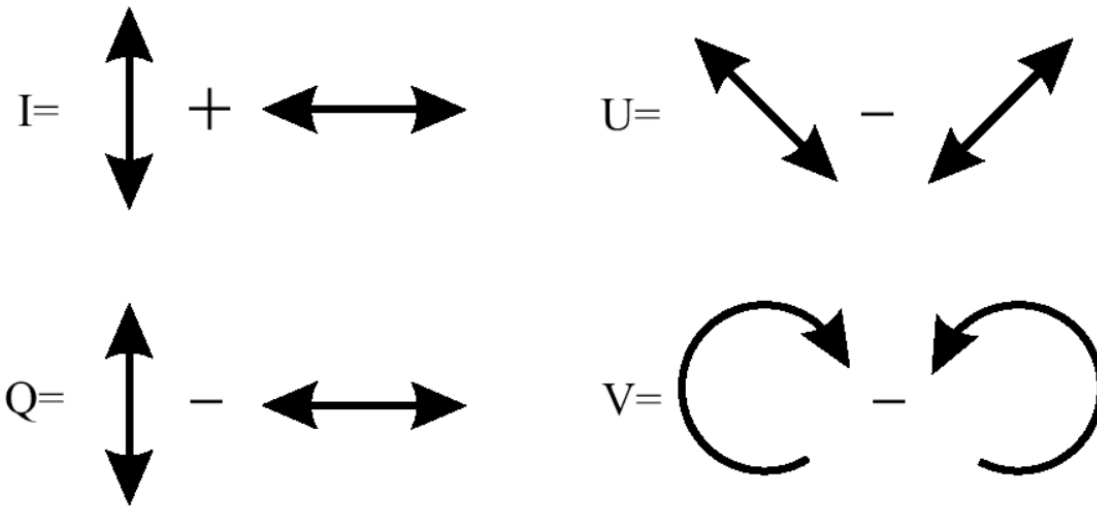
- Cannot be directly **measured**
- We measure electromagnetic radiation (**photons**)
- The measured radiation has information about the properties of the emitting plasma

Introduction



- **Polarization** is present when there is no symmetry:
 - **Scattering**: radiation pumping by anisotropic radiation
 - **Zeeman** effect: energy splitting of degenerate atomic levels
 - **Hanle** effect: relaxation of quantum coherences

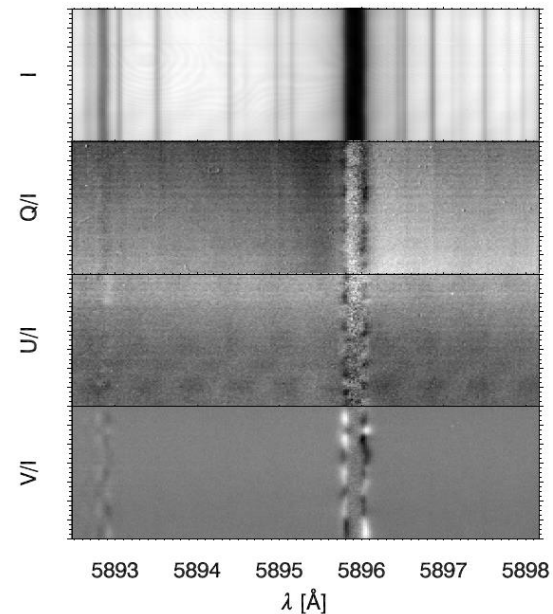
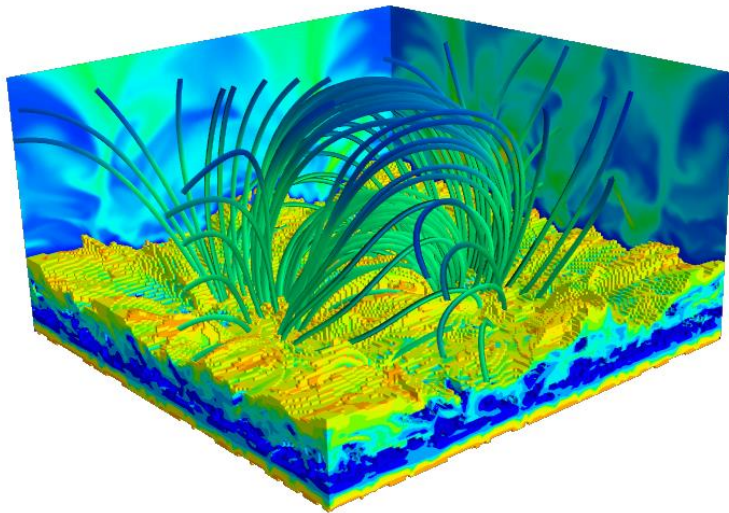
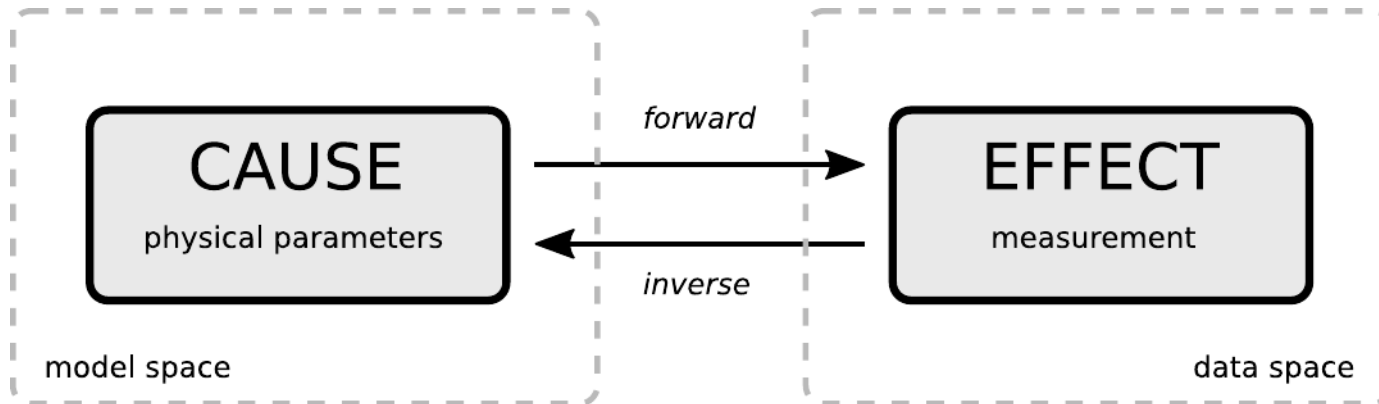
Introduction



- **Polarization** is present when there is no symmetry:
 - **Scattering**: radiation pumping by anisotropic radiation
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**Magnetic
effects**

Introduction

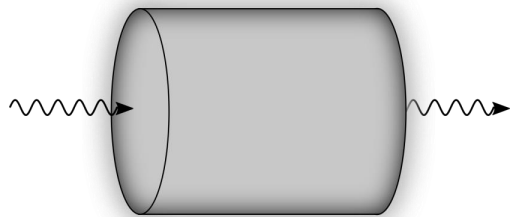


Na I D1 slit observation with ZIMPOL@IRSOL

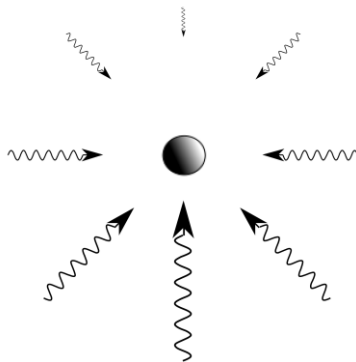
The Forward Problem Radiation Transfer

Radiation Transfer

- Describes the radiation-matter interaction
- Two parts:



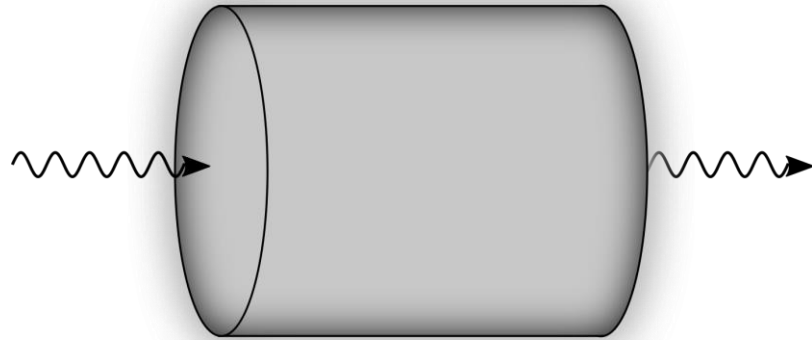
- How is the radiation propagated through a medium



- How are the atoms excited within the atmospheric radiation field

Radiation Transfer Equation

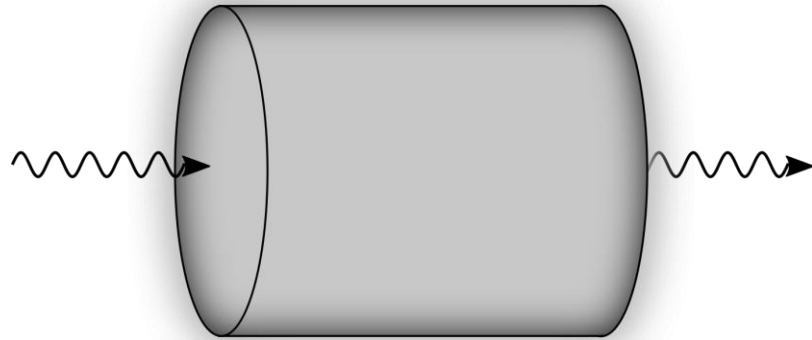
- How the radiation is modified along its propagation



$$\frac{d}{ds} \begin{pmatrix} I \\ Q \\ U \\ V \end{pmatrix} = - \begin{pmatrix} \eta_I & \eta_Q & \eta_U & \eta_V \\ \eta_Q & \eta_I & \rho_V & -\rho_U \\ \eta_U & -\rho_V & \eta_I & \rho_Q \\ \eta_V & \rho_U & -\rho_Q & \eta_I \end{pmatrix} \begin{pmatrix} I \\ Q \\ U \\ V \end{pmatrix} + \begin{pmatrix} \epsilon_I \\ \epsilon_Q \\ \epsilon_U \\ \epsilon_V \end{pmatrix}$$

Radiation Transfer Equation

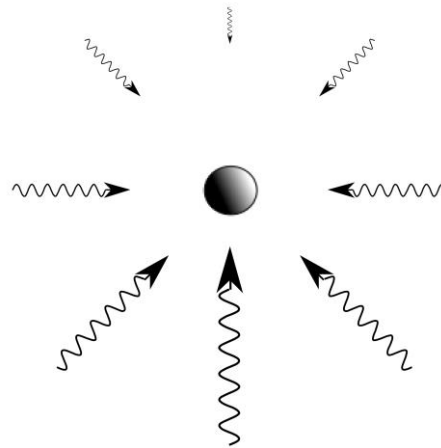
- How the radiation is modified along its propagation



$$\frac{d\mathbf{S}}{ds} = -\hat{\mathbf{K}} \cdot \mathbf{S} + \epsilon$$

Statistical Equilibrium Equations

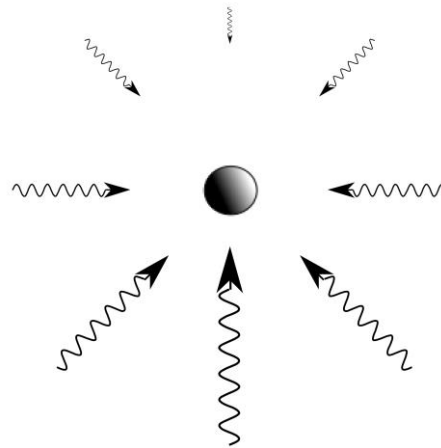
- How the atom is excited within a radiation field



$$\begin{aligned}
 \left[\frac{d}{dt} {}^L \rho_Q^K(J, J') \right]_{\text{Rad}} &= \sum_{L_\ell J_\ell J'_\ell K_\ell Q_\ell} {}^{L_\ell} \rho_{Q_\ell}^{K_\ell}(J_\ell, J'_\ell) \mathbb{T}_A(LJJ'KQ, L_\ell J_\ell J'_\ell K_\ell Q_\ell) \\
 &+ \sum_{L_u J_u J'_u K_u Q_u} {}^{L_u} \rho_{Q_u}^{K_u}(J_u, J'_u) [\mathbb{T}_E(LJJ'KQ, L_u J_u J'_u K_u Q_u) + \mathbb{T}_S(LJJ'KQ, L_u J_u J'_u K_u Q_u)] \\
 &- \sum_{J'' J''' K' Q'} {}^L \rho_{Q'}^{K'}(J'', J''') [\mathbb{R}_E(LJJ'KQJ''J'''K'Q') + \mathbb{R}_S(LJJ'KQJ''J'''K'Q') + \\
 &\quad \mathbb{R}_A(LJJ'KQJ''J'''K'Q')]
 \end{aligned}$$

Statistical Equilibrium Equations

- How the atom is excited within a radiation field



$$\frac{d\rho_i}{dt} = \sum_{j \neq i} R_{ji} \rho_j - \sum_j R_{ij} \rho_i = 0$$

Radiation Transfer

$$\frac{d\mathbf{S}}{ds} = -\hat{\mathbf{K}} \cdot \mathbf{S} + \epsilon$$

$$\frac{d\rho_i}{dt} = \sum_{j \neq i} R_{ji} \rho_j - \sum_j R_{ij} \rho_i = 0$$

Radiation Transfer

$$\frac{d\mathbf{S}}{ds} = -\hat{\mathbf{K}}(\rho) \mathbf{S} + \epsilon(\rho, \mathbf{S})$$

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Radiation Transfer

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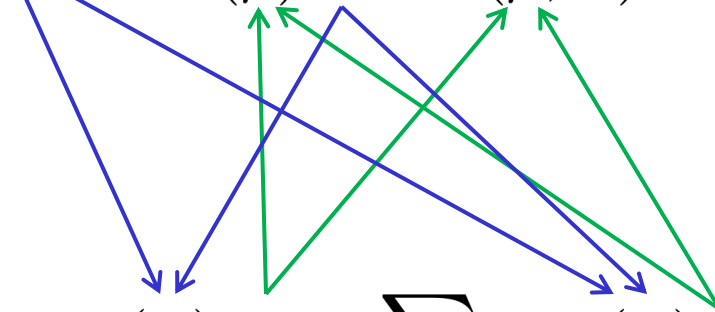
The diagram illustrates the relationship between the radiation transfer equation and the rate of change of population density. The first equation, $\frac{d\mathbf{S}}{ds} = -\hat{\mathbf{K}}(\rho) \mathbf{S} + \epsilon(\rho, \mathbf{S})$, describes the change in the radiation field \mathbf{S} along a path s . The second equation, $\frac{d\rho_i}{dt} = \sum_{j \neq i} R_{ji}(\mathbf{S}) \rho_j - \sum_j R_{ij}(\mathbf{S}) \rho_i = 0$, describes the rate of change of the population density ρ_i for state i . Blue arrows indicate the flow from the terms in the first equation to the terms in the second equation: from $\frac{d\mathbf{S}}{ds}$ to $\frac{d\rho_i}{dt}$, from $\hat{\mathbf{K}}(\rho)$ to $R_{ij}(\mathbf{S})$, and from $\epsilon(\rho, \mathbf{S})$ to $R_{ji}(\mathbf{S})$. Green arrows indicate the flow from the terms in the second equation back to the terms in the first equation: from ρ_j to $\hat{\mathbf{K}}(\rho)$, from ρ_i to $\hat{\mathbf{K}}(\rho)$, and from $R_{ji}(\mathbf{S})$ to $\epsilon(\rho, \mathbf{S})$.

Radiation Transfer

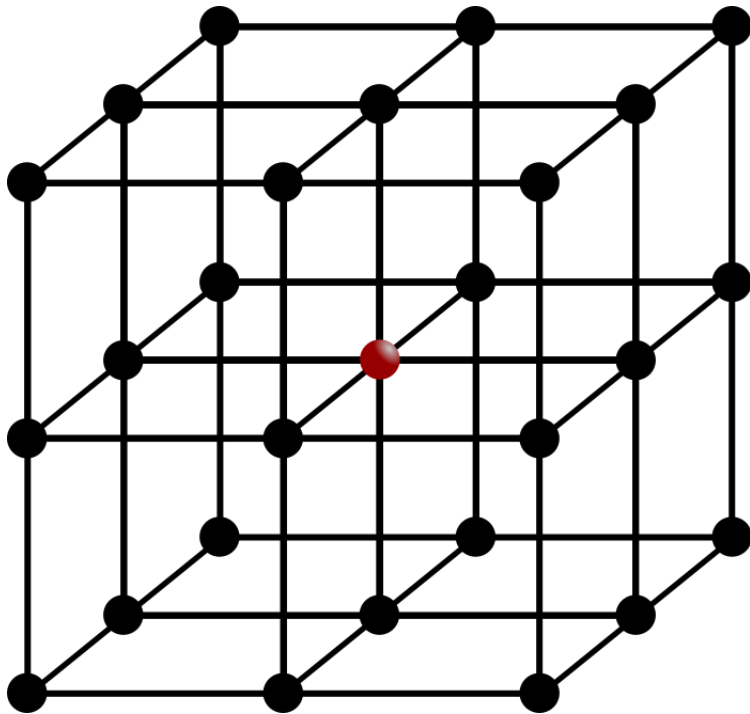
- Coupled
- Non-linear
- Non-local
- Most difficult (costly) problem to solve in solar physics

$$\frac{d\mathbf{S}}{ds} = -\hat{\mathbf{K}}(\rho) \mathbf{S} + \epsilon(\rho, \mathbf{S})$$

$$\frac{d\rho_i}{dt} = \sum_{j \neq i} R_{ji}(\mathbf{S}) \rho_j - \sum_j R_{ij}(\mathbf{S}) \rho_i = 0$$

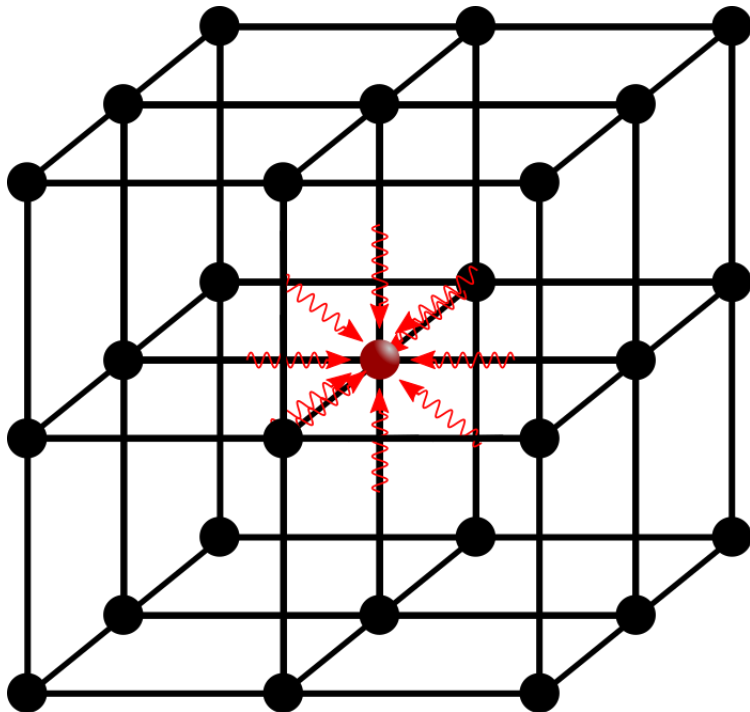


Radiation Transfer



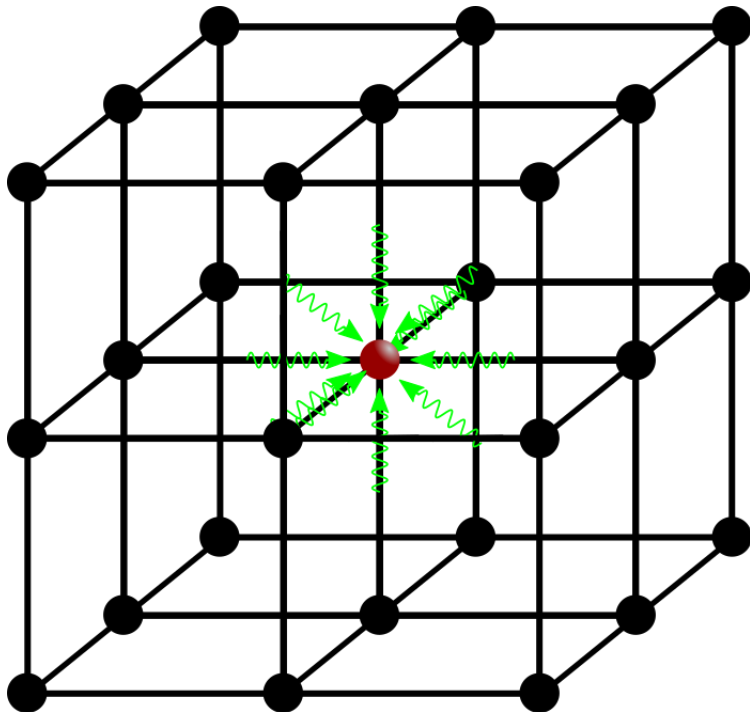
- Spatial nodes:
 $n_x, n_y, n_z \sim 10^2$
 $n_x \cdot n_y \cdot n_z \sim 10^6$

Radiation Transfer



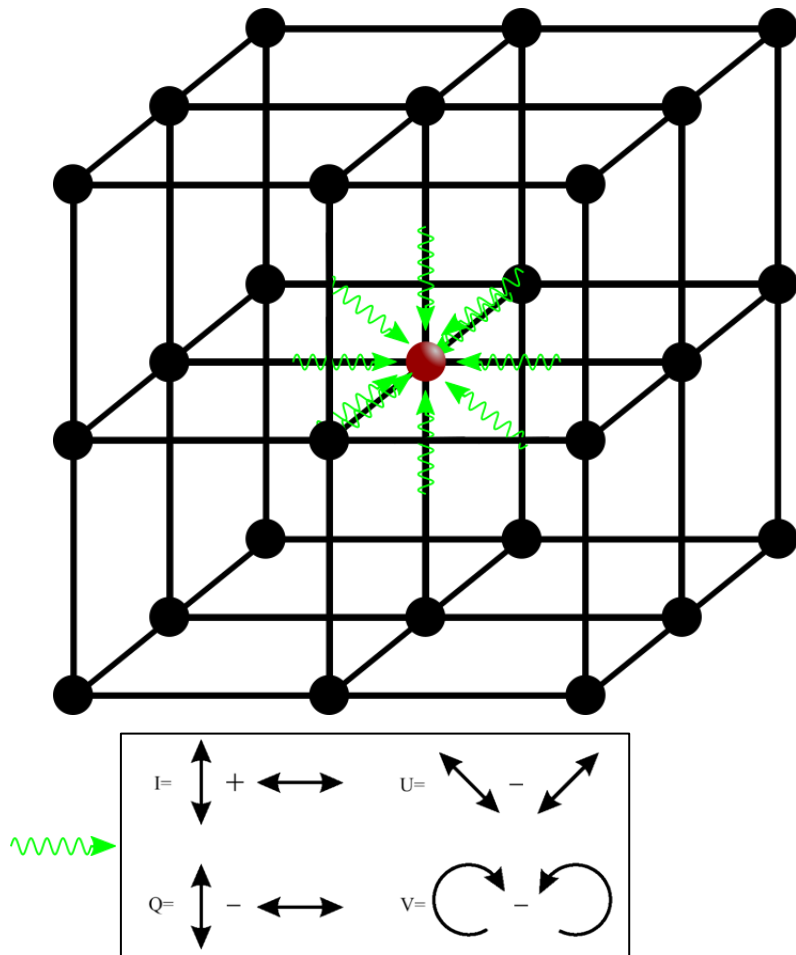
- Spatial nodes:
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- Directions:
 $n_\Omega \sim 10^2$

Radiation Transfer



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 $n_x, n_y, n_z \sim 10^2$
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- Directions:
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- Frequencies (per line):
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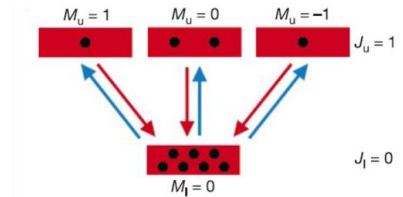
Radiation Transfer



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- Directions:
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- Frequencies (per line):
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- Polarization:
 $n_s \sim 4$

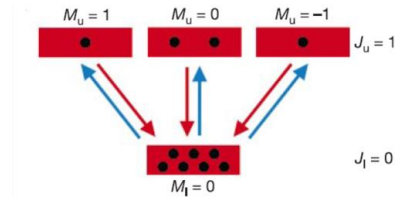
Radiation Transfer

- Simplest problem: two-level atom $J_l = 0$; $J_u = 1$



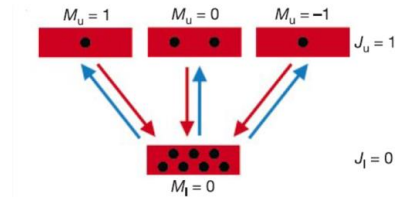
Radiation Transfer

- Simplest problem: two-level atom $J_l = 0; J_u = 1$
- Statistical equilibrium:
 - 10 unknowns per spatial node $\rightarrow 10^7$



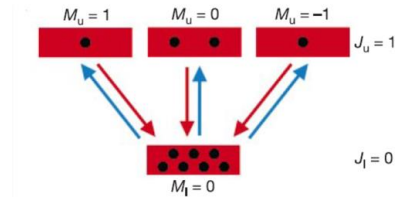
Radiation Transfer

- Simplest problem: two-level atom $J_l = 0; J_u = 1$
- Statistical equilibrium:
 - 10 unknowns per spatial node $\rightarrow 10^7$
- Radiation field:
 - Polarization in every node, frequency, and direction, 10^{10} unknowns



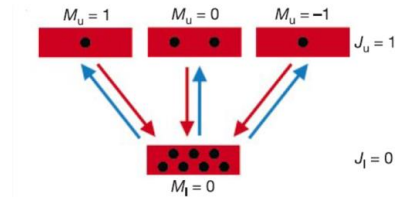
Radiation Transfer

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- The problem is iterative \rightarrow repeat $\sim 10^2$ times



Radiation Transfer

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- Statistical equilibrium:
 - 10 unknowns per spatial node $\rightarrow 10^7$
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 - Polarization in every node, frequency, and direction, 10^{10} unknowns
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Parallelization is a **must**

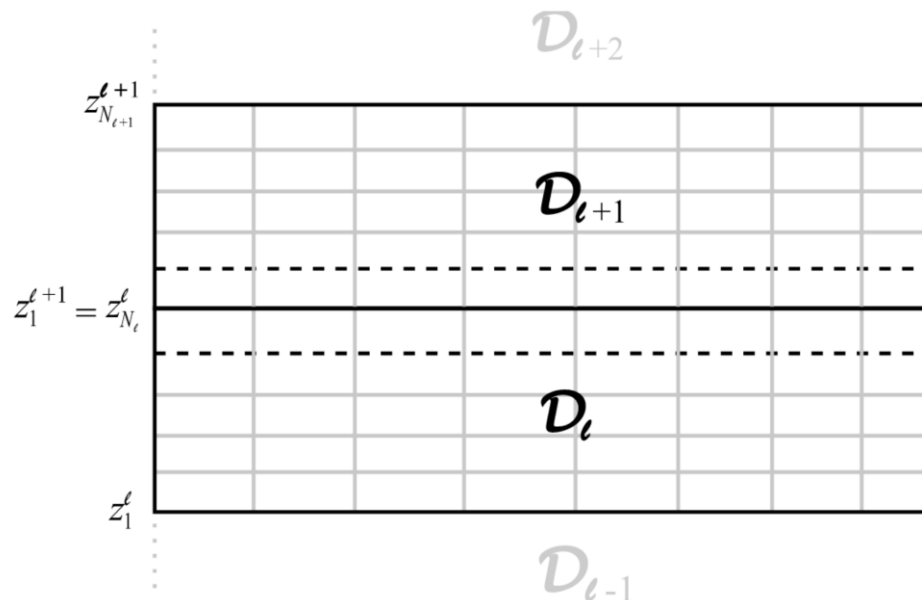
PORTA

PORTA

- **Library** to solve the problem of the generation and transfer of polarized radiation in 3D atmospheres
Štěpán & Trujillo Bueno (2013)
- **Modules** to solve specific problems
- Almost **linear** scaling with #CPU

PORTA

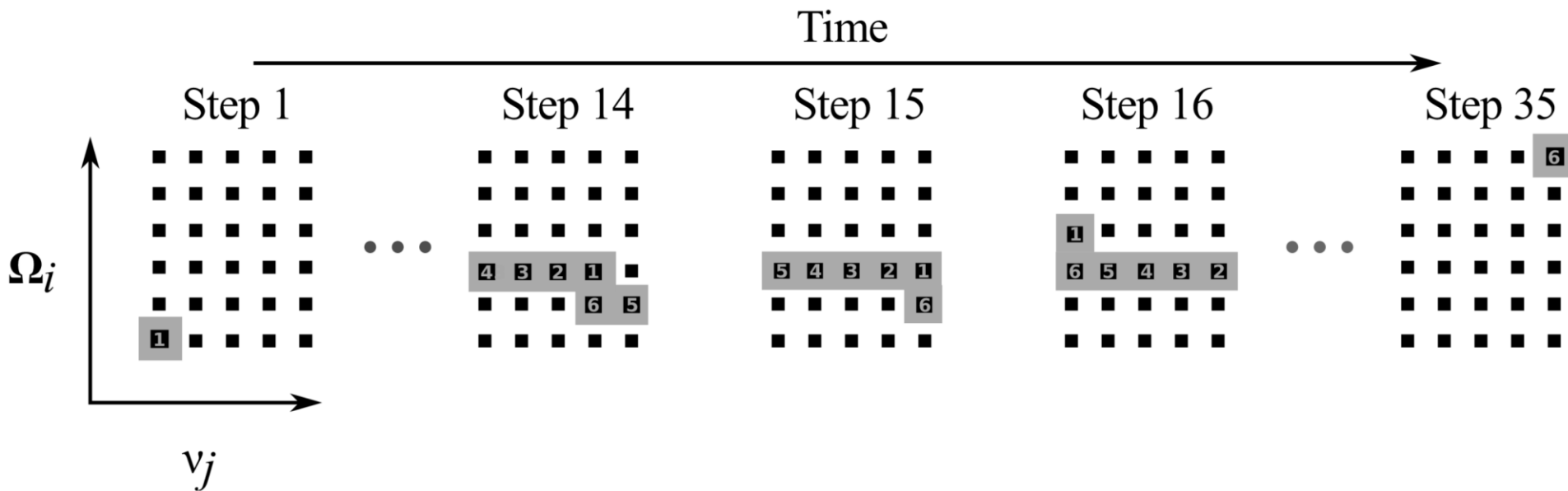
- Domain decomposition:
 - Distributes work
 - Eases memory constraints



Štěpán & Trujillo Bueno (2013)

PORTA

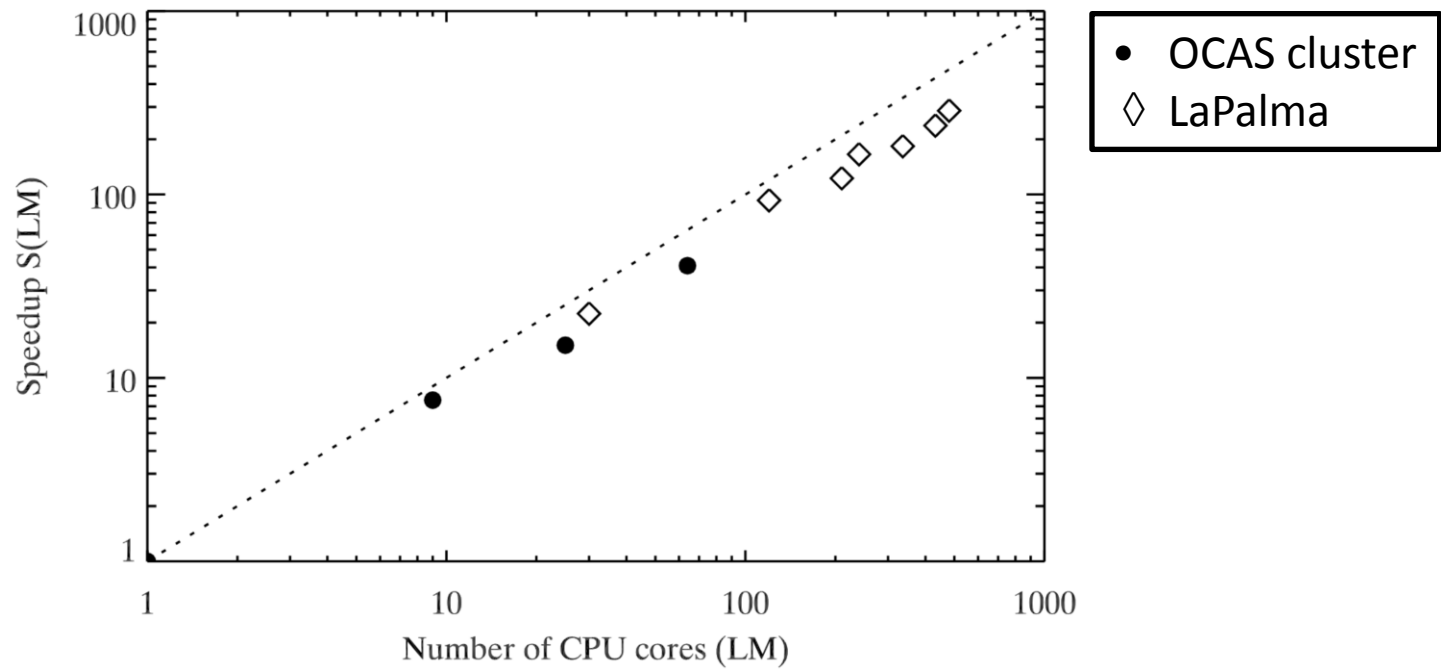
- Snake algorithm:



Štěpán & Trujillo Bueno (2013)

PORTA

- Linear scaling:



Štěpán & Trujillo Bueno (2013)

PORTA

- Close to be public with the modules:
 - Coherent scattering (Jaume Bestard, J., del Pino Alemán, T., and Štěpán, J.)
 - General two-level (Štěpán, J.)
 - General multi-level (del Pino Alemán, T.)

Applications in MareNostrum

PORTA

- Every application uses the PORTA code
- But every application is different
 - Preliminar investigations (1D theoretical studies)
 - Preliminar computations (preparation of 3D models and computation of intermediate quantities)
 - Different modules
- I will only talk about the very final results of some investigations
 - And only about one or two results of the chosen ones

PORTA

- Hydrogen Lyman- α :
 - Theoretical study (Štěpán et al. (2015))
 - Diagnostics of CLASP data (Trujillo Bueno et al. (2018))
- Hydrogen Balmer- α :
 - Theoretical study (Jaume Bestard PhD. thesis, WIP)
- Calcium 4227 Å:
 - Theoretical study (Jaume Bestard PhD. thesis, WIP)
- Calcium H-K and infrared triplet:
 - Theoretical study (Štěpán and Trujillo Bueno (2016))
 - Comparison with observations (Jurčák et al. (2018))

PORTA

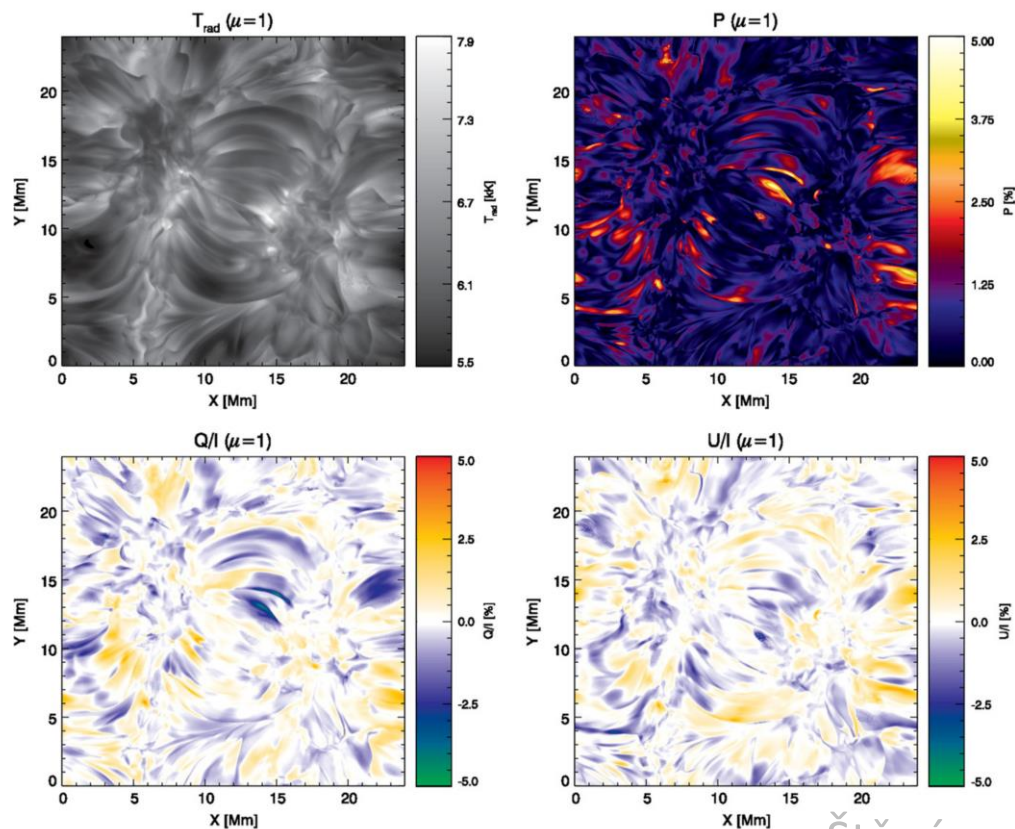
- Mg II k-line:
 - Theoretical study (del Pino Alemán PhD. thesis (2015))
- Sr I 4607 Å:
 - Theoretical study and comparison with observations (del Pino Alemán et al. (2018))
- Radiation transfer theoretical study:
 - Polarization with horizontal inhomogeneities (Tichý et al. (2015))

Some Results

Diagnostic of CLASP observations

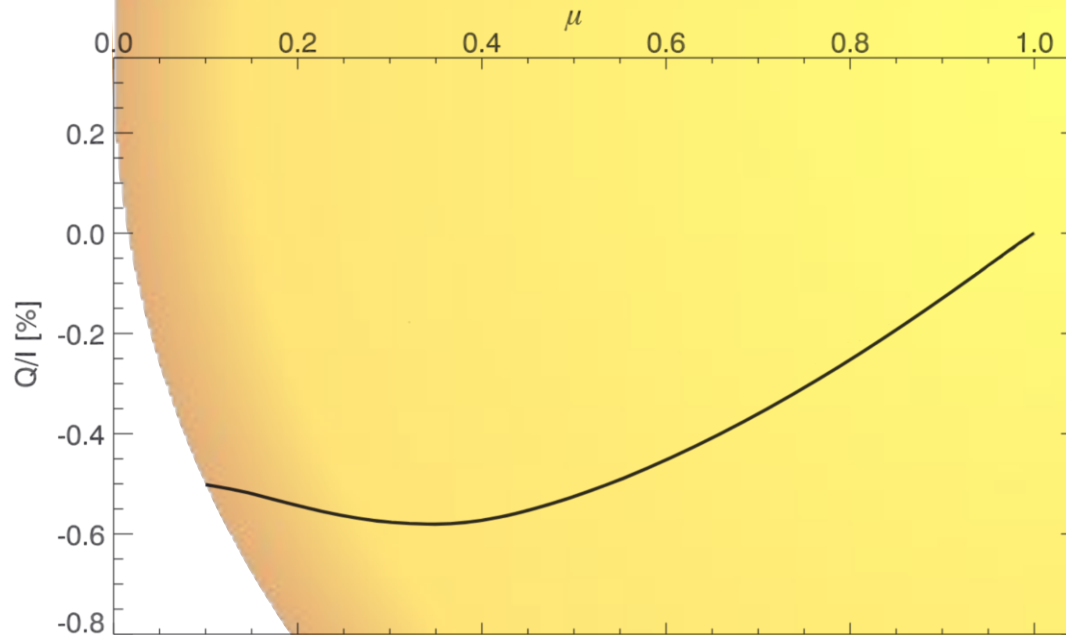
Diagnostic of CLASP observations

- We carried out the detailed radiation transfer modeling of the Lyman- α line



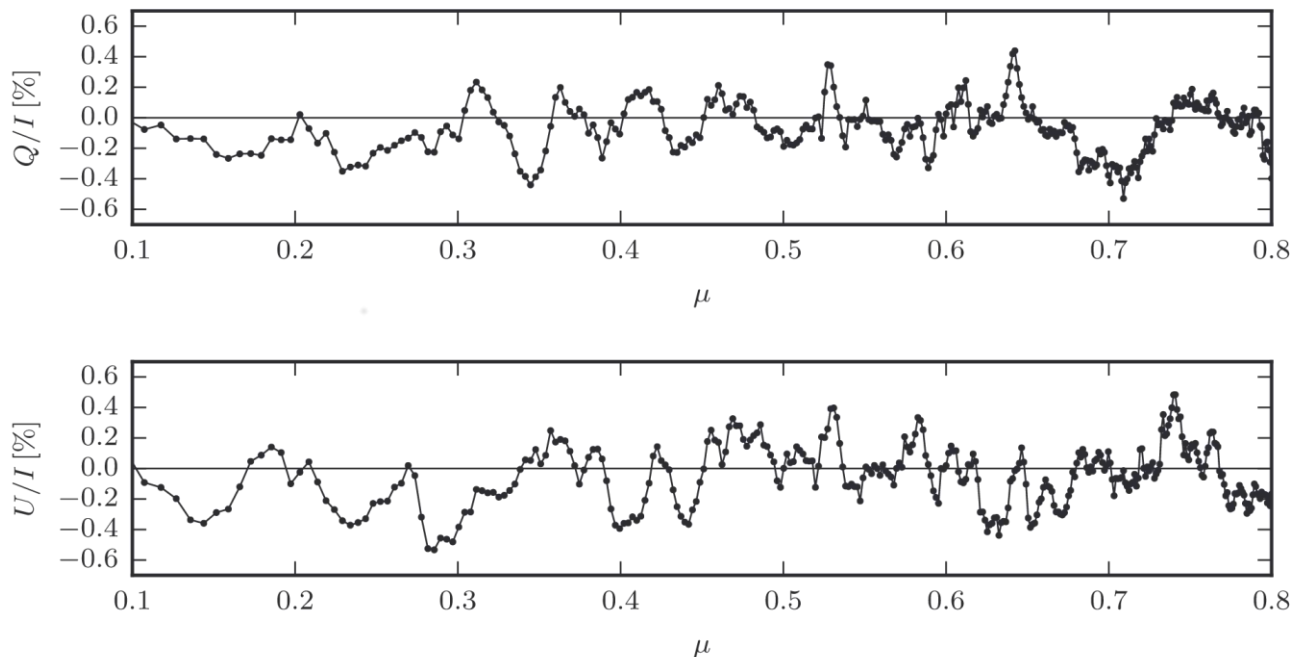
Diagnostic of CLASP observations

- Theory predicted center-to-limb variation of the linear polarization in the center of the Lyman- α line



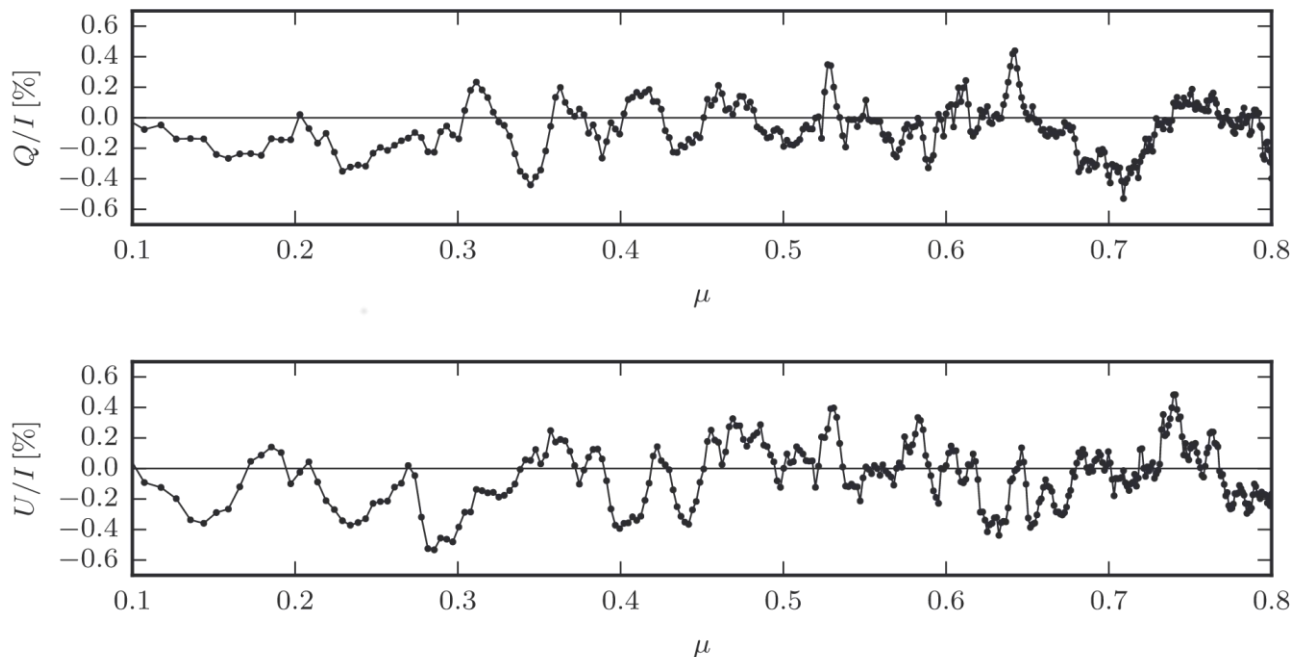
Diagnostic of CLASP observations

- But observations did not show any



Diagnostic of CLASP observations

- But observations did not show any
- The 3D model could not reproduce the observations



Diagnostic of CLASP observations

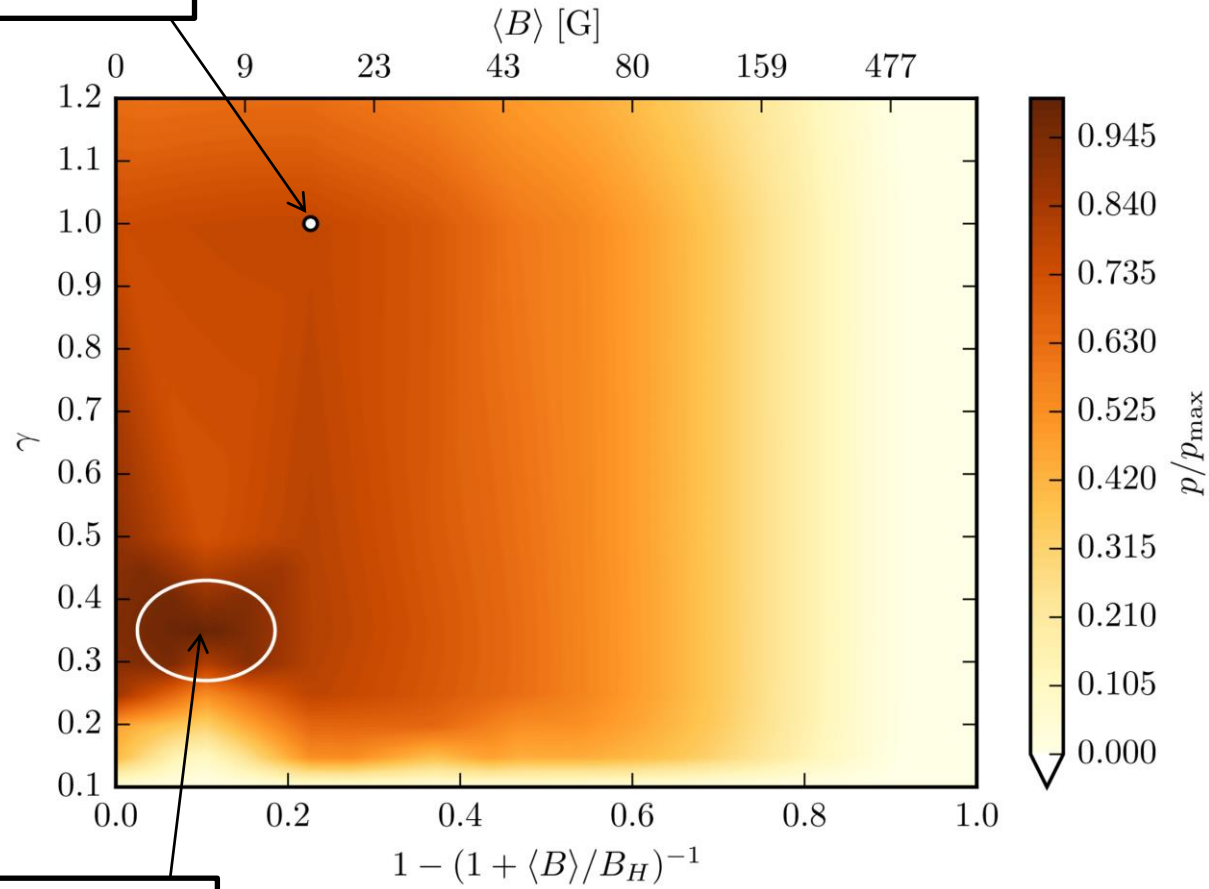
- What other properties should the 3D model have?
- We introduced two parameters:
 - Geometrical complexity (compression factor)
 - Magnetization (magnetic field strength factor)
- Bayesian approach: what is the combination of parameters with the most likelihood?

Diagnostic of CLASP observations

Original 3D model

Observations indicate:

- More corrugation
- Less magnetization



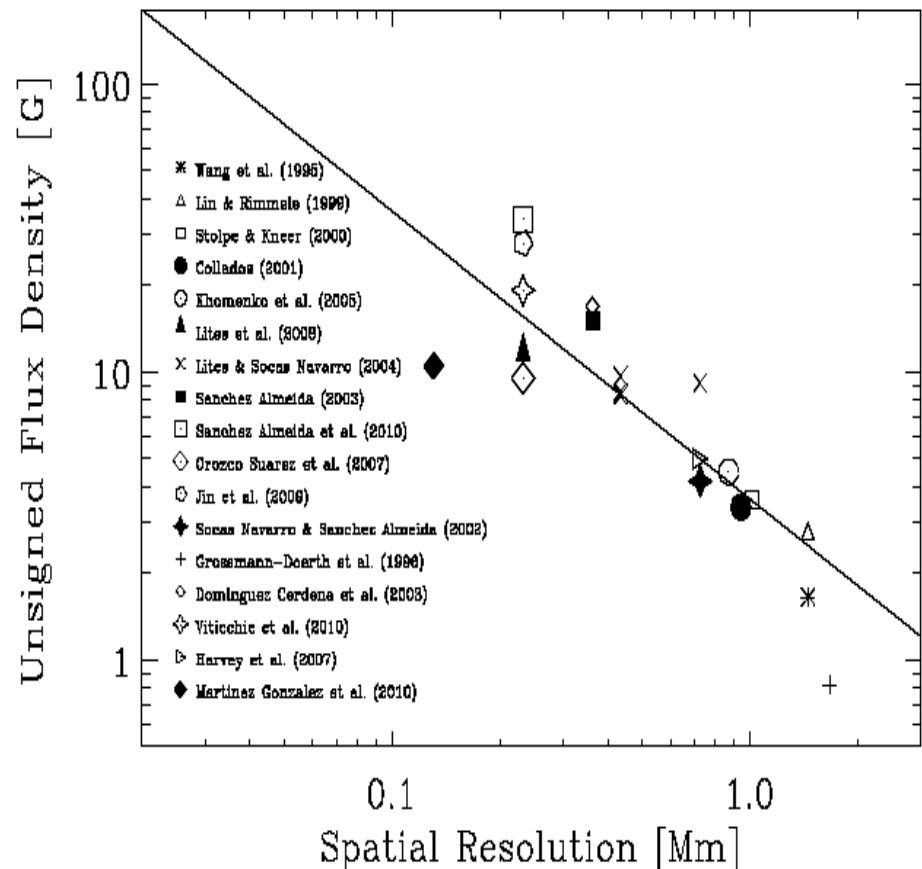
Maximum likelihood

Trujillo Bueno et al. (2018)

Photospheric small
scale magnetization

Photospheric small scale magnetism

- The photospheric magnetic field has structure at very small scales
- The resolution limit of solar observations limits our detection capability
- More field the better the resolution

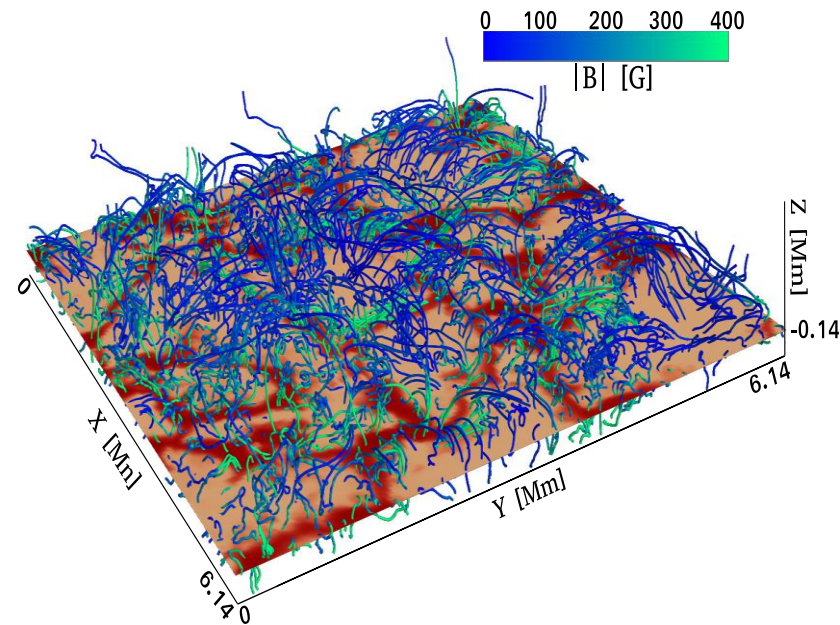


Photospheric small scale magnetism

- Several authors with different approaches have provided estimations of $\langle B \rangle$.
 - Stenflo (1982): $>10\text{G}$
 - Faurobert et al. (1995): $10\text{-}20\text{G}$
 - Trujillo Bueno et al. (2004): 100G

Photospheric small scale magnetism

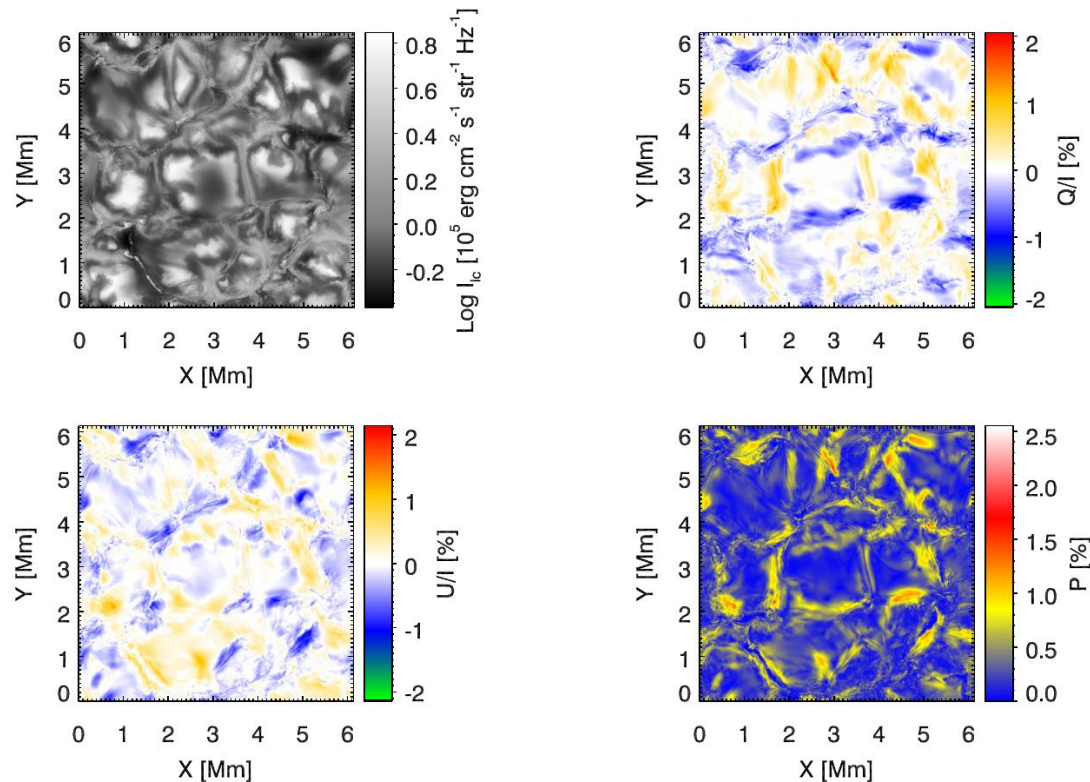
- All previous 3D (Hanle) calculations assumed non spatially resolved magnetic fields
- Rempel (2014) provided 3D magneto-convection simulations with significant level of small-scale magnetic activity
 $\langle B \rangle \approx 170\text{G}$ at height $\approx 0\text{km}$



del Pino Alemán et al. (2018)

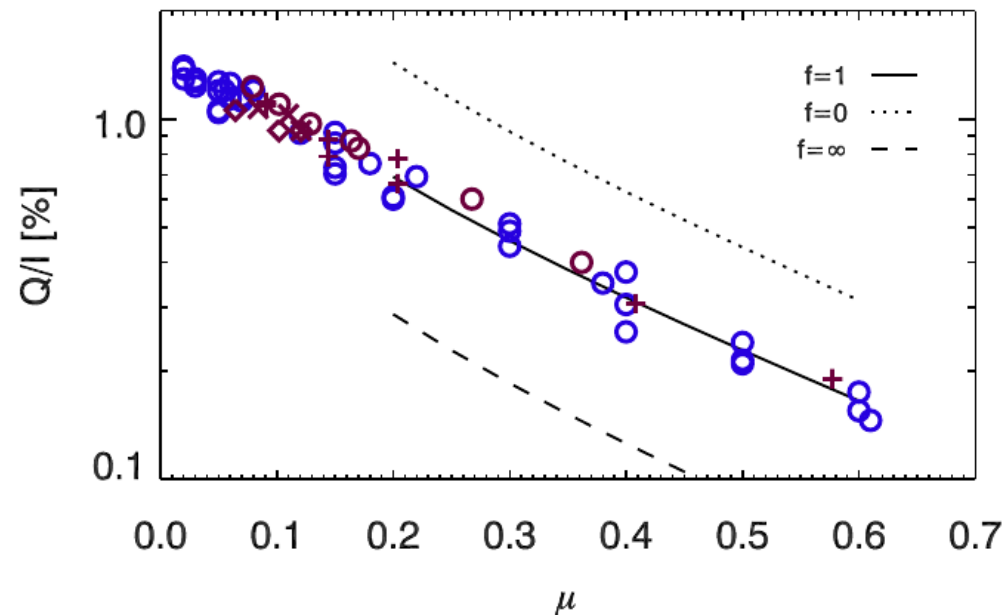
Photospheric small scale magnetism

- We carried out the detailed radiation transfer modeling of Sr I 4607Å in this model



Photospheric small scale magnetism

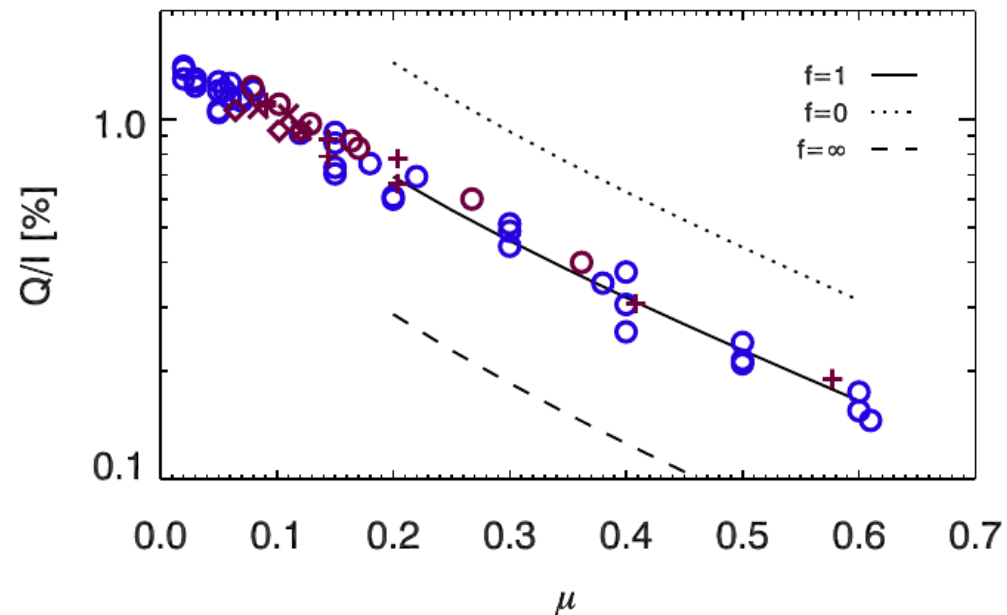
- We carried out the detailed radiation transfer modeling of Sr I 4607Å in this model
- Compare with center-to-limb variation observations



del Pino Alemán et al. (2018)

Photospheric small scale magnetism

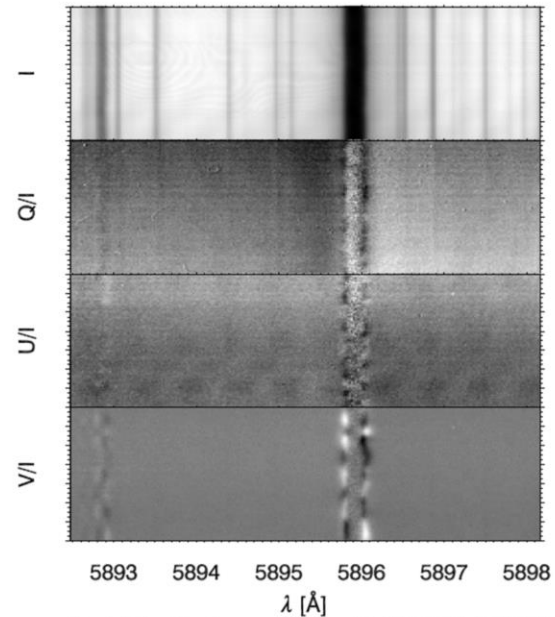
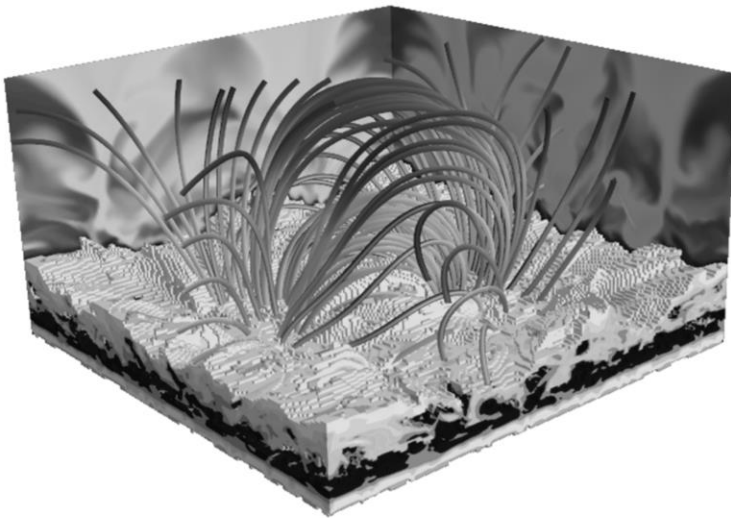
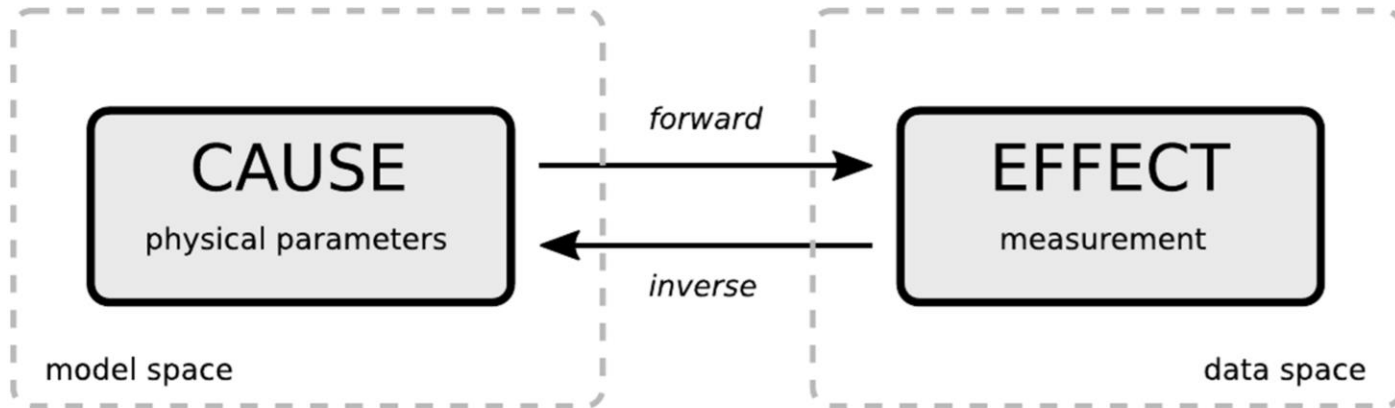
- We carried out the detailed radiation transfer modeling of Sr I 4607Å in this model
- Compare with center-to-limb variation observations
- The level of magnetization is compatible with the observations:
 $\langle B \rangle \approx 170\text{G}$ at the surface
- (and a lot of other theoretical/statistical considerations I have no time to talk about now)



del Pino Alemán et al. (2018)

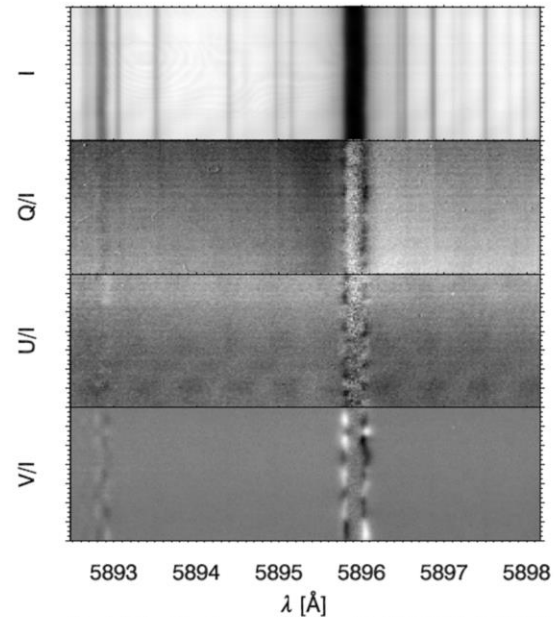
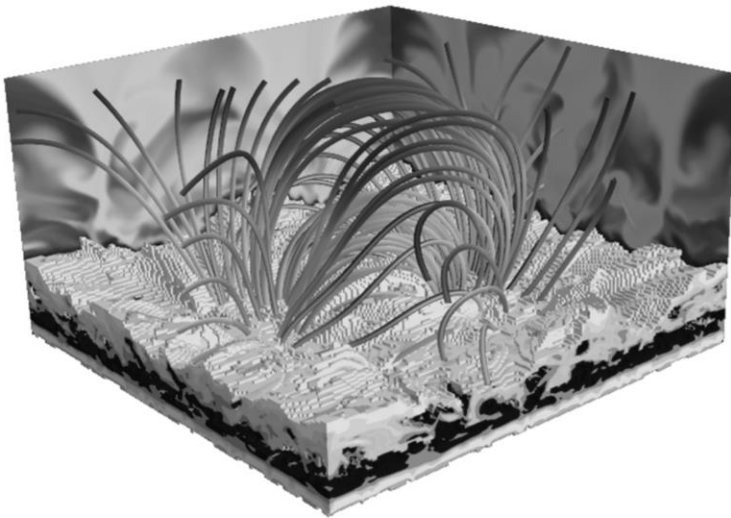
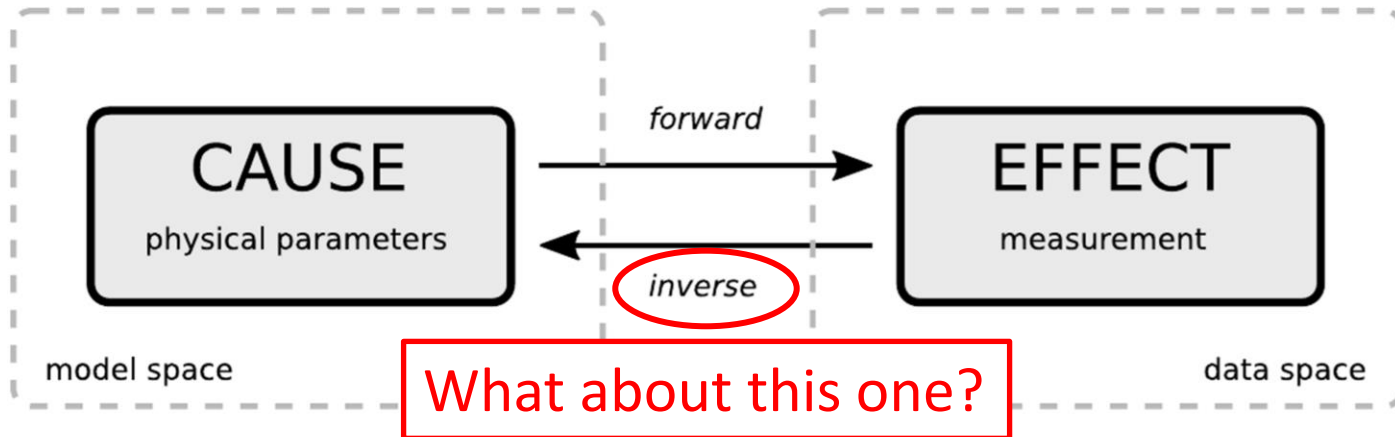
That is a lot of
forward modeling
but...

Introduction



Na I D1 slit observation with ZIMPOL@IRSOL

Introduction



Na I D1 slit observation with ZIMPOL@IRSOL

The inversion problem

The inversion problem

- Find the x -parameters that fulfills

$$y = F(x) + \varepsilon$$

y : data

F : forward problem

ε : noise

The inversion problem

- Find the x -parameters that fulfills

$$y = F(x) + \varepsilon$$

y : data
 F : forward problem
 ε : noise

- We cannot do $x = F^{-1}(y)$. We solve the optimization problem

$$x = \underset{x}{\operatorname{argmin}} \|y - F(x)\|_2^2$$

The inversion problem

- Find the x -parameters that fulfills

$$y = F(x) + \varepsilon$$

y : data
 F : forward problem
 ε : noise

- We cannot do $x = F^{-1}(y)$. We solve the optimization problem

$$x = \underset{x}{\operatorname{argmin}} \|y - F(x)\|_2^2$$

- Still ill-posed. We introduce some regularization

$$x = \underset{x}{\operatorname{argmin}} \{ \|y - F(x)\|_2^2 + g(x) \}$$

- Example: sparsity $g(x) = \lambda \|x\|_0$; best subset

The inversion problem

- We want to find the simplest model that reproduce the data
- Evaluating F the minimum amount of times
- With a method that scales linearly with #CPU
- First attempt (to my knowledge) of inversions with a 3D forward solver

The inversion problem

Picture by Rob Glover

- Sparsity is a rare occurrence in the '*real*' space



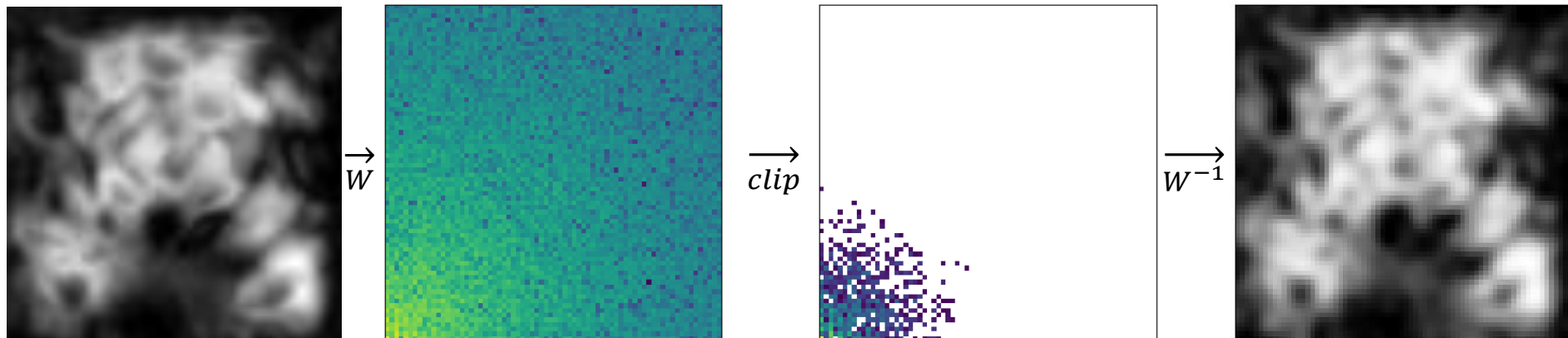
The inversion problem

Picture by Rob Glover

- Sparsity is a rare occurrence in the '*real*' space



- But can be a common occurrence in a transformed space



The inversion problem

- Find the q -parameters, images of the x -parameters that fulfills

$$q \subseteq DCT(x): F(IDCT(q)) + \varepsilon = y$$

The inversion problem

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- But we cannot test every q -parameter to check if it is relevant

The inversion problem

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$$q \subseteq DCT(x): F(IDCT(q)) + \varepsilon = y$$

- But we cannot test every q -parameter to check if it is relevant
- Time constrains force us to keep only the smoother modes

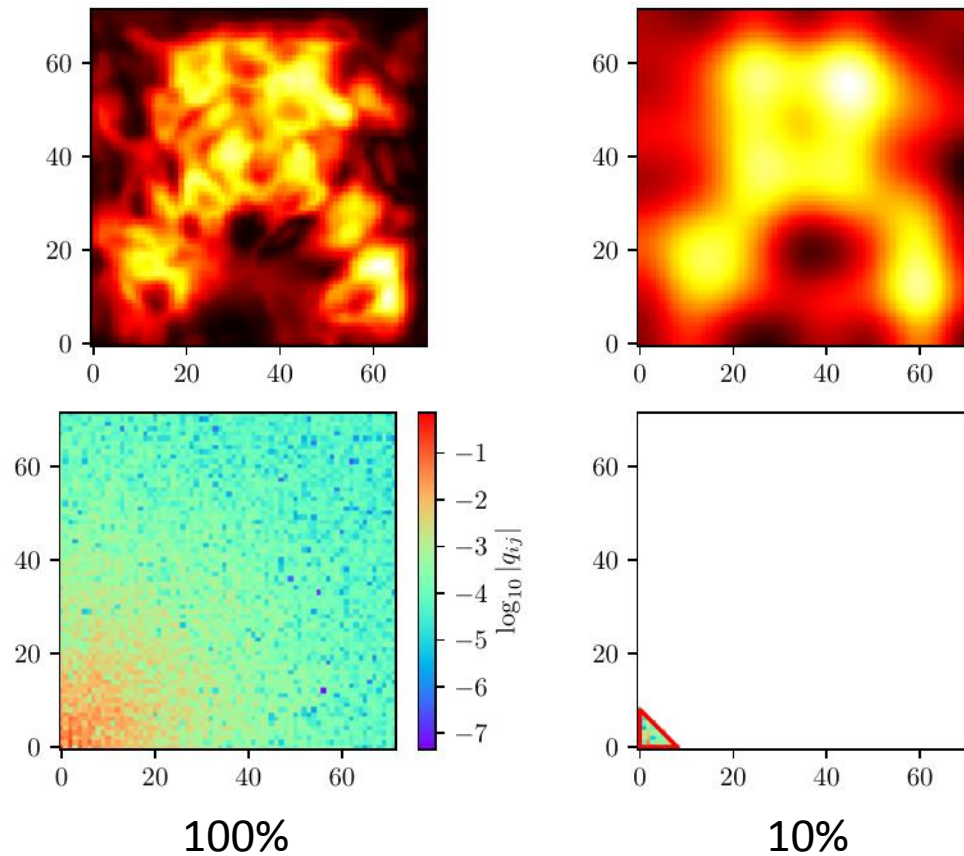
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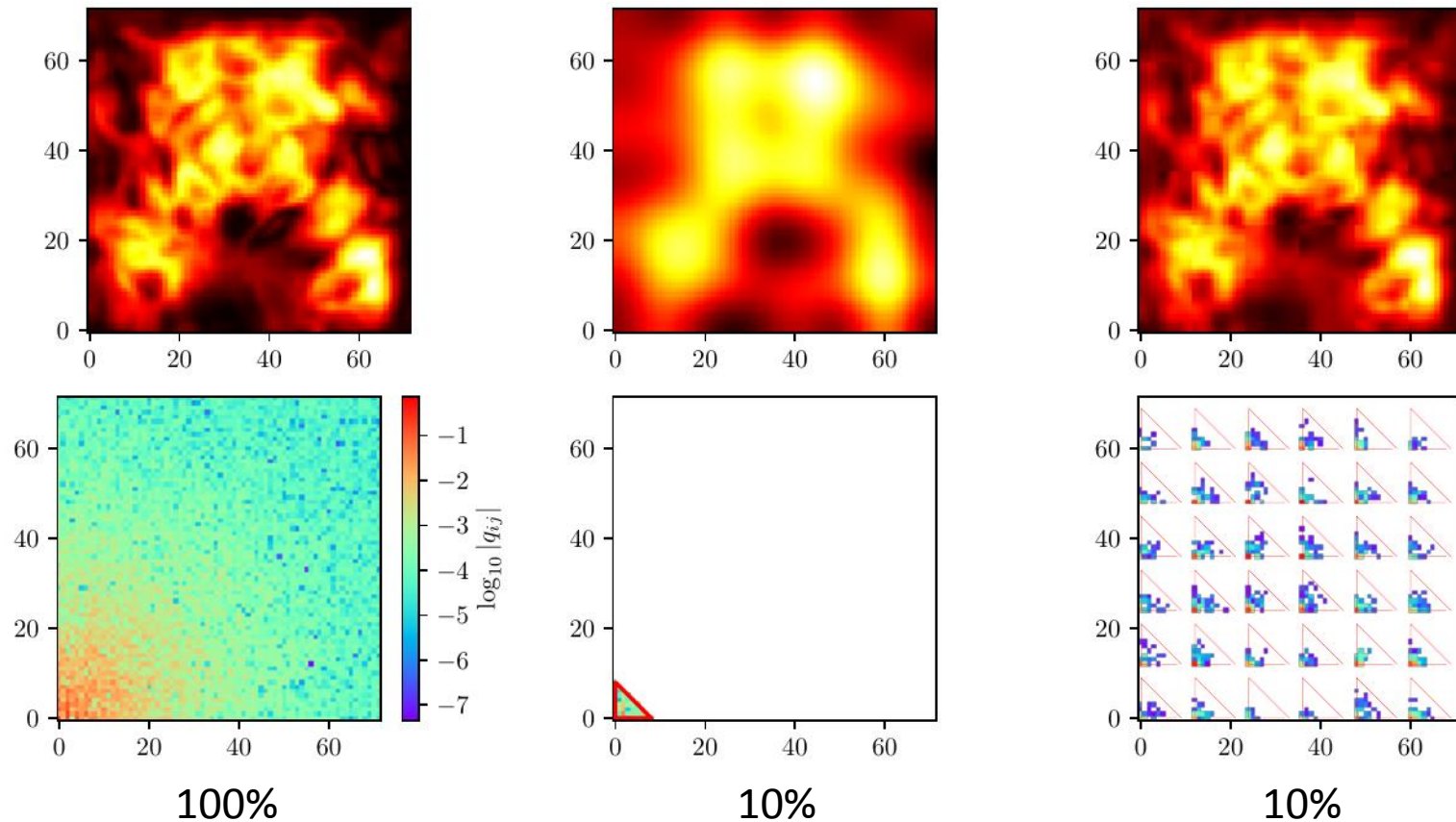
- But we cannot test every q -parameter to check if it is relevant
- Time constrains force us to keep only the smoother modes
- But we loose details

The inversion problem



The inversion problem

- We introduce tiling



The inversion problem

- We loose a bit of spatial coherency

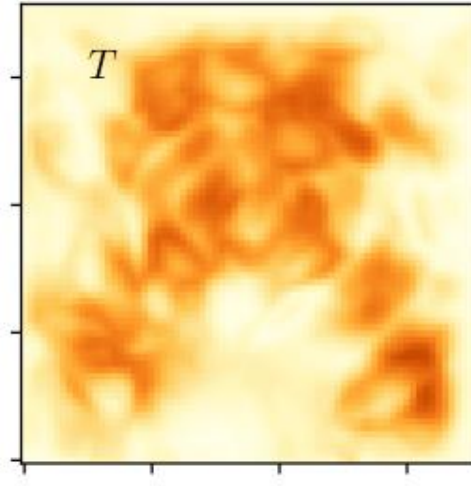
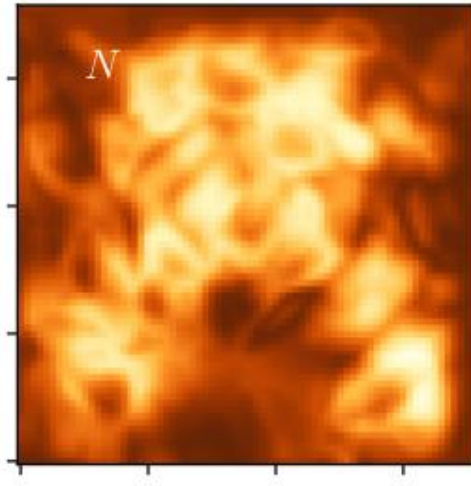
The inversion problem

- We loose a bit of spatial coherency
- But everything is still consistent because the tiles interact radiatively in the forward solver
- We win:
 - More resolution capability
 - Parallel computation of each mode in every tile speed-up by 2-3 orders of magnitude

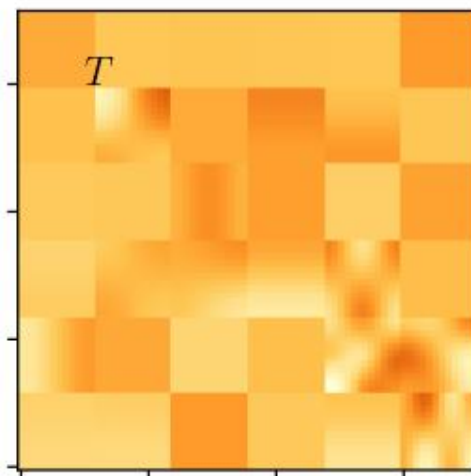
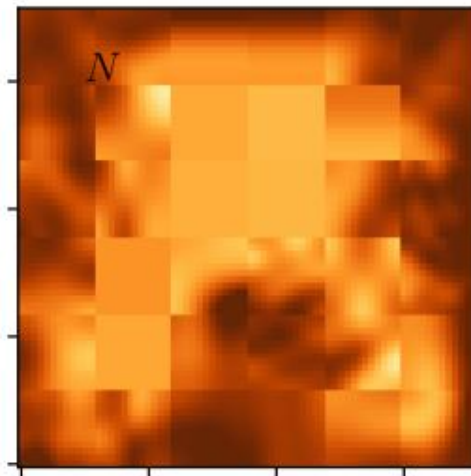
The inversion problem

- Algorithm:
 - start with very sparse solution
 - iterate until convergence
 - if the agreement not good enough:
increase modes and repeat
- Result:
 - model with the minimum number of parameters
 - physically consistent

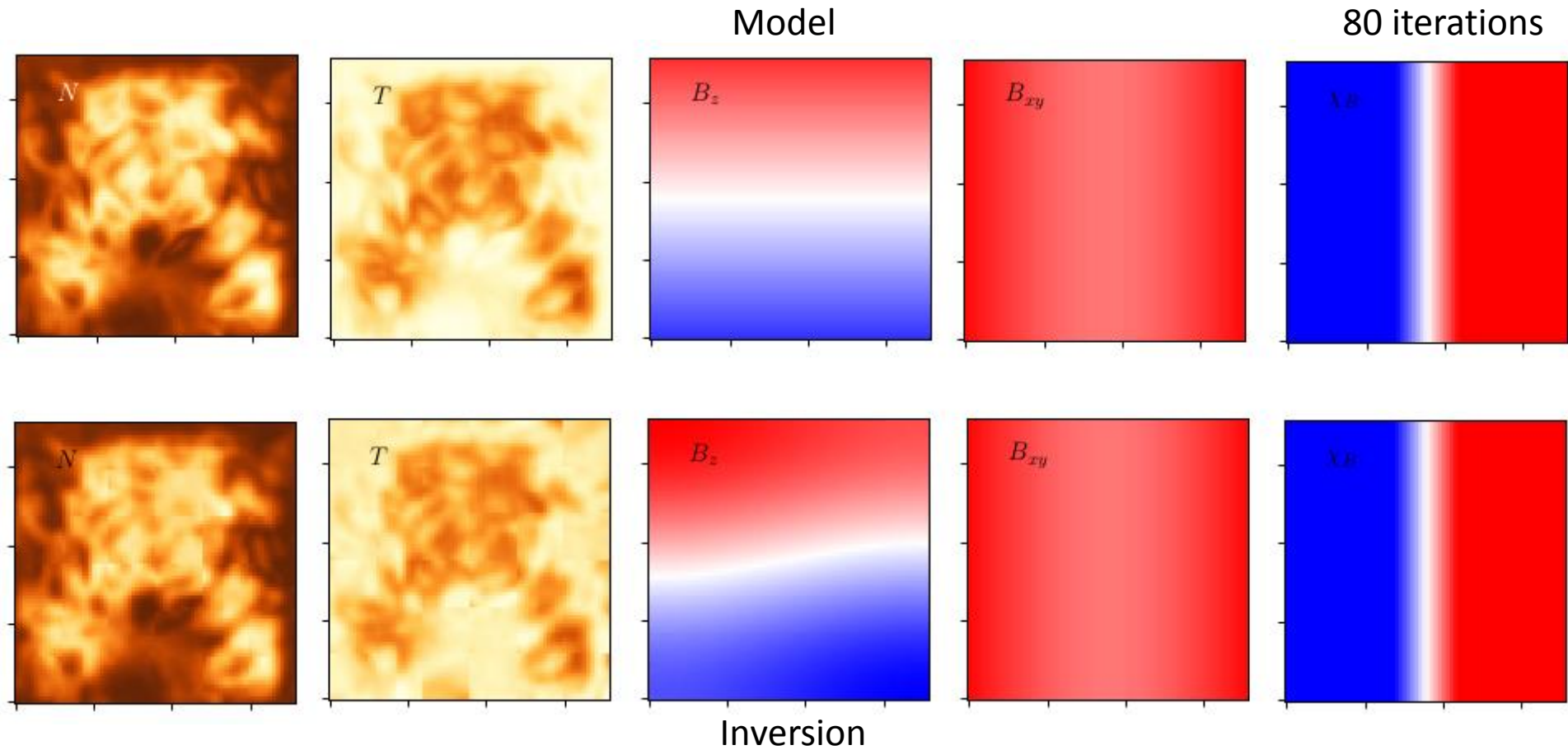
The inversion problem



30 iterations



The inversion problem



The inversion problem

- This is a recent test of feasibility
- We are planing to apply this inversion tool to real data very soon

Thank you for your attention