



"SuperStretch"  
PICS collaboration  
Lyon-Madrid  
in progress

# Ice nucleation:

What we have learned so far, and  
what still needs to be understood

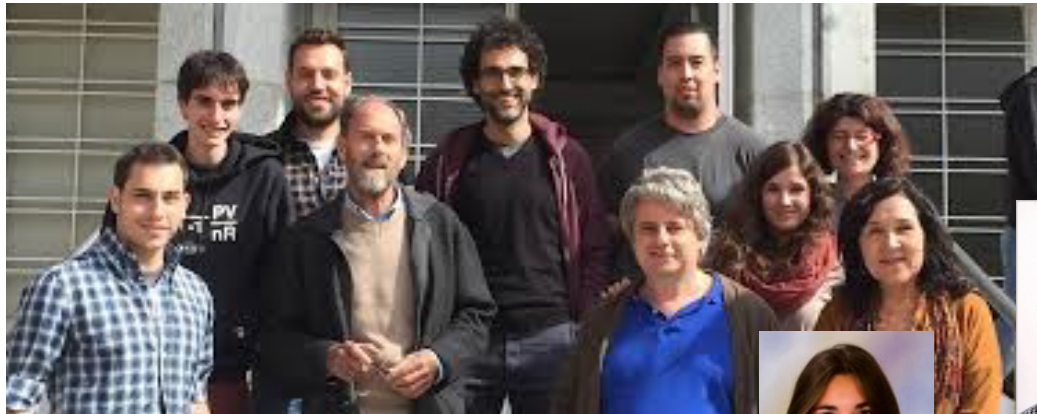
Chantal Valeriani

Universidad Complutense de Madrid  
RES workshop, Santiago, September 2017

# The ice nucleation team

## Universidad Complutense de Madrid

- Eduardo Sanz
- Jose Luis Abascal
- Carlos Vega
- Jorge Espinosa
- Pablo Rosales
- Alberto Zaragoza
- Pablo Montero
- Anibal Pandiani
- Victor Cruces
- Marta Garcia
- Raul Martinez
- Miguel Angel Gonzalez
- Francisco Alarcon



## Universidad Politecnica de Madrid

- Jorge Ramirez

## Universidad de Guanajuato (Mexico)

- Analaura Benavides



# Ice nucleation



However, nucleation is still not fully understood in many systems

Why is the mechanism of water freezing still not known?

Experimentally: it is not possible to directly observe the formation of ice clusters being too small and too short lived

**SIMULATIONS CAN PLAY A RELEVANT ROLE!**



# Outline

- How to simulate water in a computer
- How to study nucleation in a computer
- Ice nucleation
- Ice nucleation at high pressure
- Ice nucleation from salty water
- Future work

# Scientific impact of our work on nucleation, thank to RES

Since 2013, when we started the nucleation project

We have published 18 papers (among those 2 PNAS, 1 PRL, 1 JACS and 1 J.Phys.Chem.Lett.)

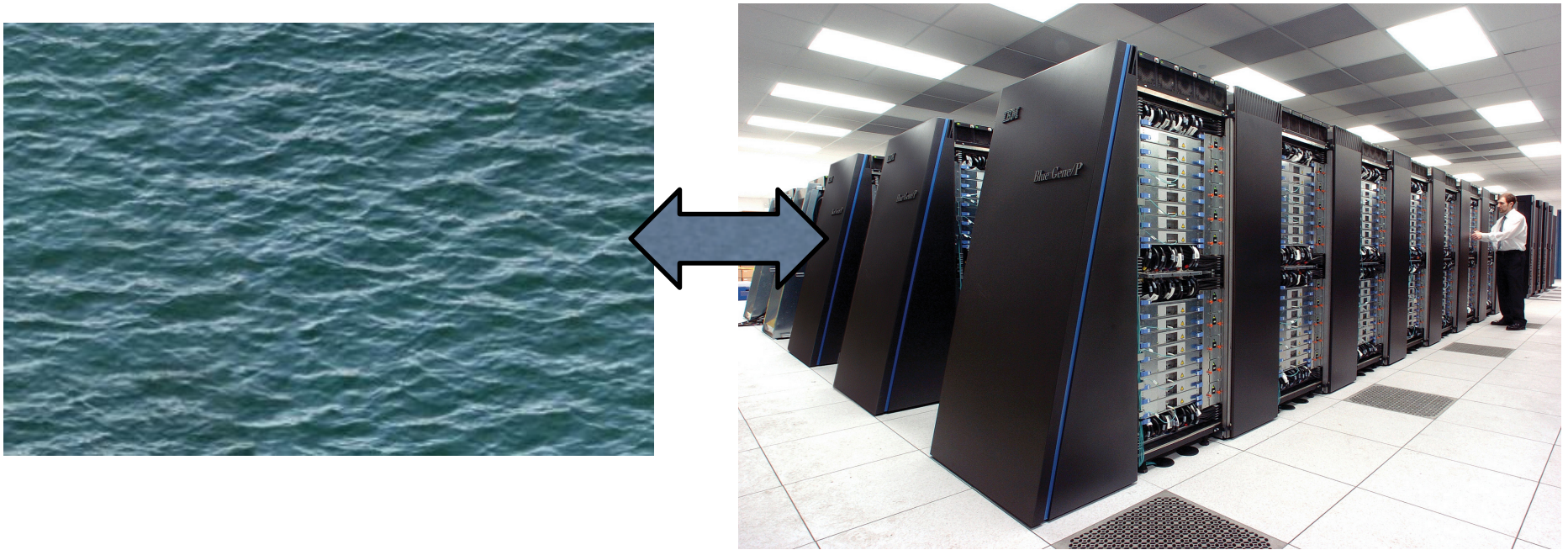
Our papers on nucleation have already received 299 citations (WOS, September 2017)

We have been invited to internationally renown conferences to present our work (such as LIQUIDS, Stat Phys, CECAM, Thermodynamics)

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# Simulations: water in a computer



how do we simulate water at a microscopic level?

# How can we study water in a computer?

Classical statistical mechanics

$$-\nabla_{\vec{R}_i} (E_e(\vec{R}^N) + V_N) = m_i \frac{d^2 \vec{R}_i}{dt^2}$$

$$E = \frac{\int (E_e(\vec{R}^N) + V_N) e^{-\beta(E_e(\vec{R}^N) + V_N)} d\vec{R}^N}{\int e^{-\beta(E_e(\vec{R}^N) + V_N)} d\vec{R}^N}$$

- (1) Empirical expression for  $E_e(\vec{R}^N) + V_N$
- +
- (2) Classical statistical mechanics



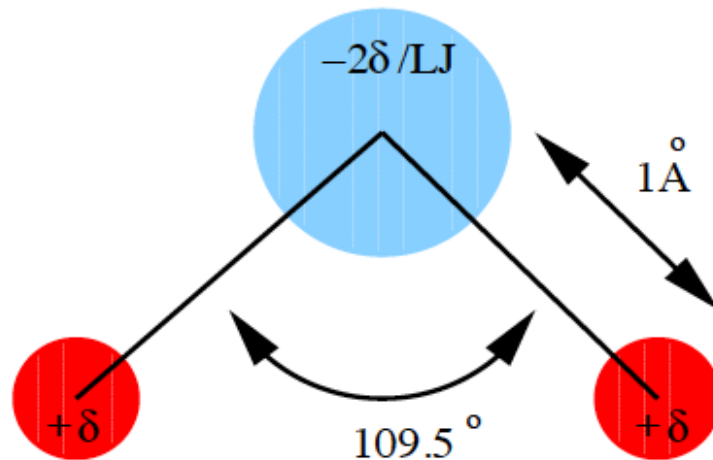
# Which water models?

SPC/E, Berendsen et al.

(1987)

TIP3P, Jorgensen et al. (1983)

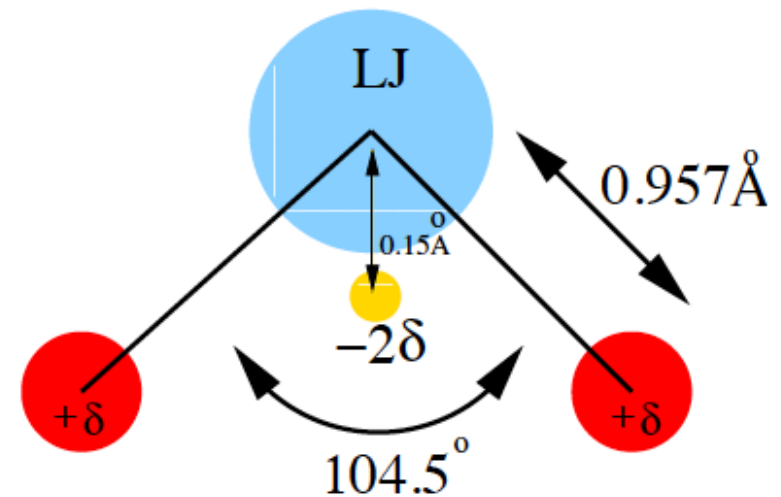
TIP4P, Jorgensen et al. 1983



1 center LJ

3 charges

SPC/E = 3000 citas



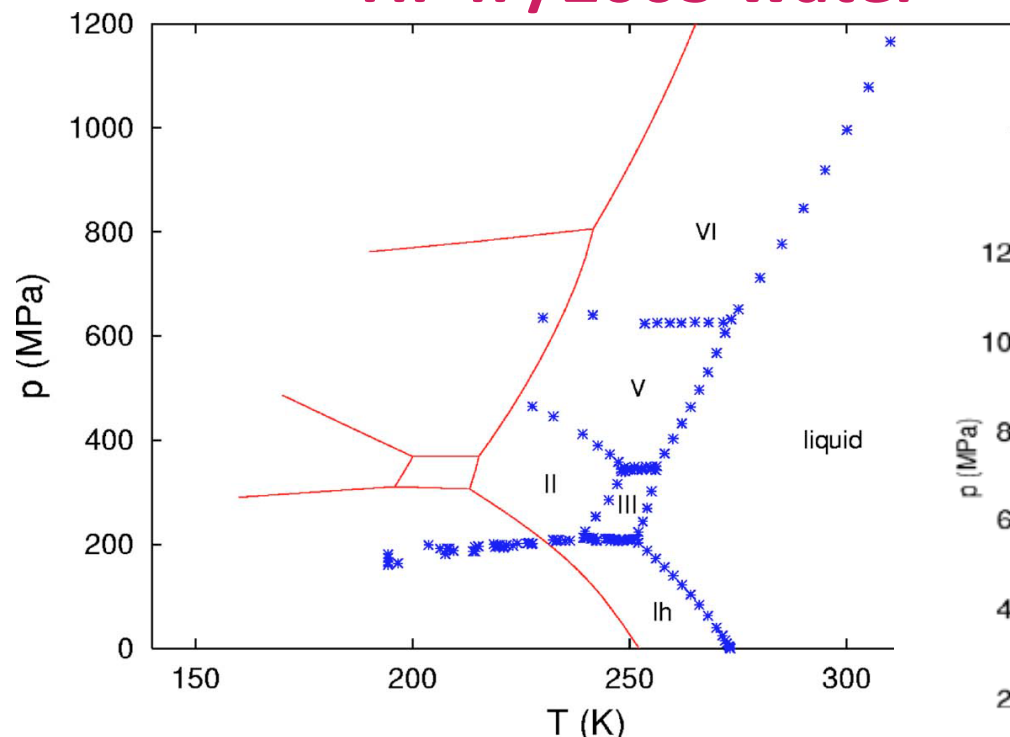
1 center LJ

3 charges

TIP3P+TIP4P = 10000 citas

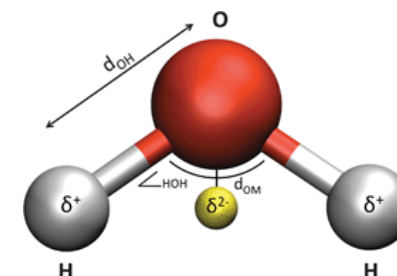
# Which water models?

## TIP4P/2005 water

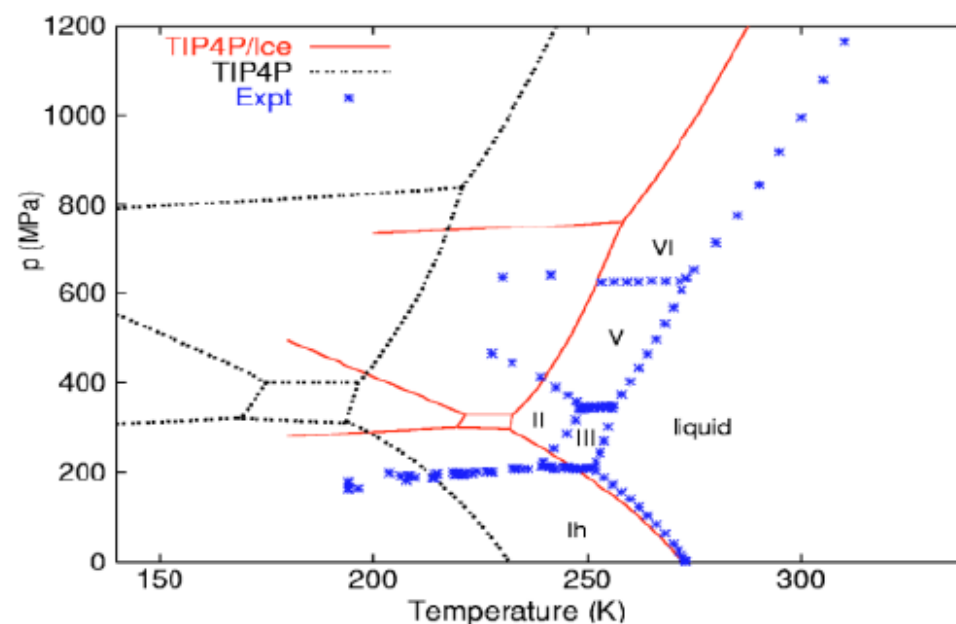


Abascal and Vega, JCP 123 234505

Abascal, Sanz, Vega, PCCP 11 556 (2009)



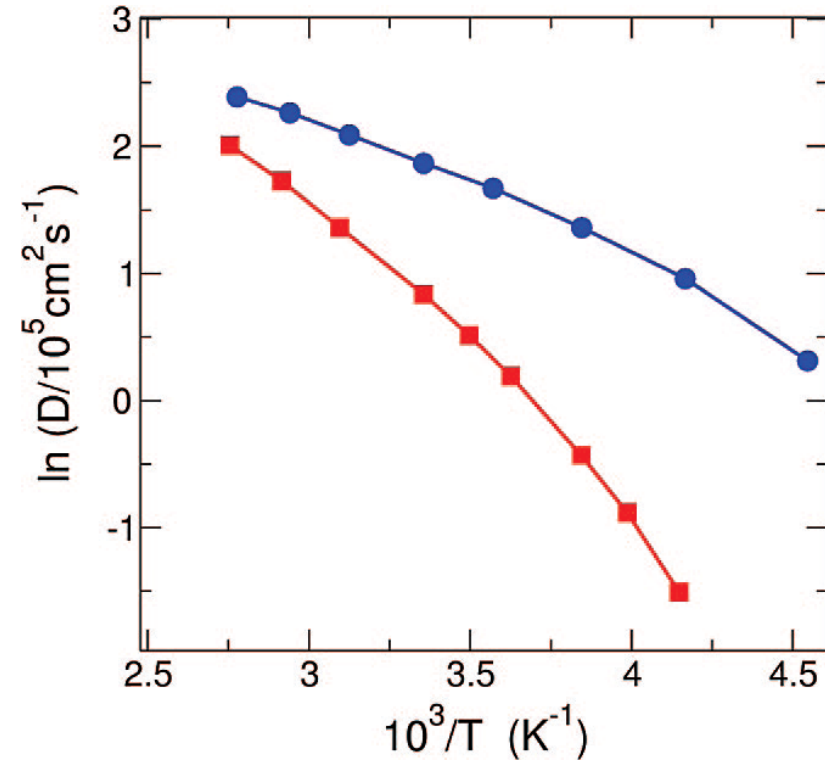
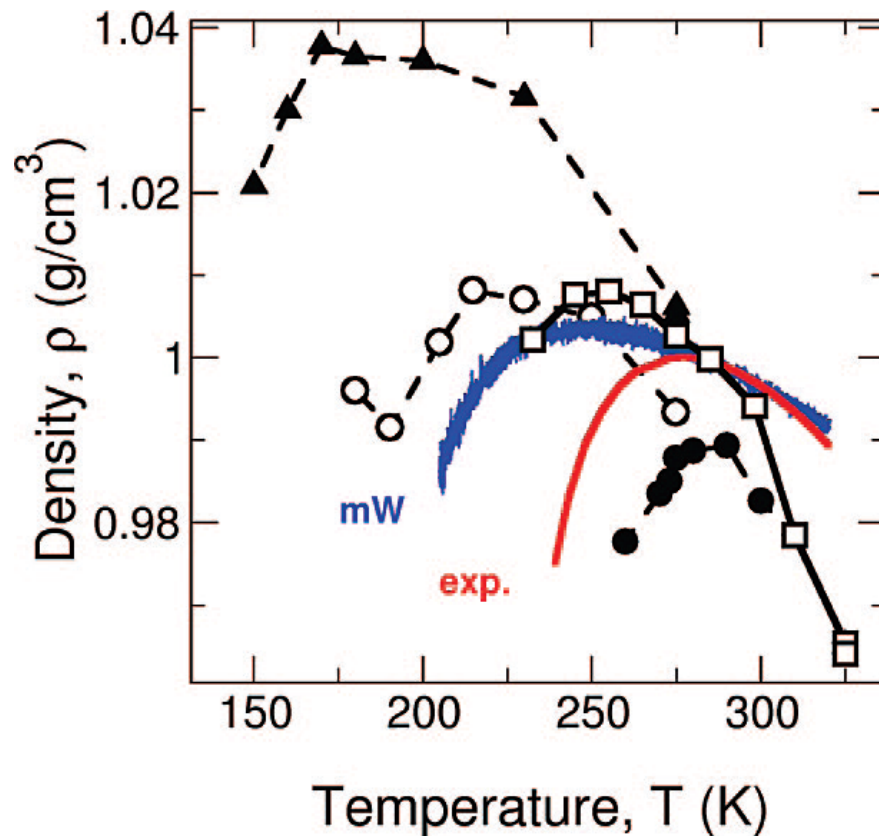
## TIP4P/ICE water



J.L.F. Abascal, E.Sanz, R.Garcia  
Fernandez and C.Vega,  
J. Chem.Phys. 122 234511 (2005)

# Which water models?

## mW water



Molinero&Moore, J.Phys.Chem.B(2009)

- ✓coarse grained model
- ✓tetrahedral environment mimicked by a 3-body angular potential
- ✓it reproduces ice melting temperature

# Outline

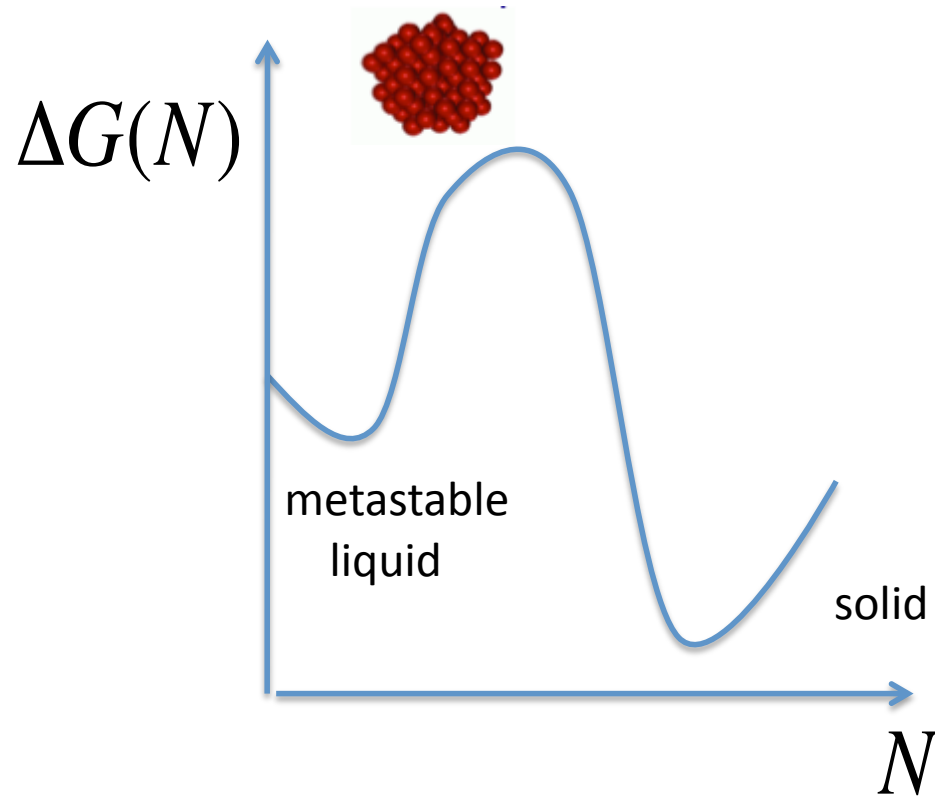
- How to simulate water in a computer
- How to study nucleation in a computer:  
the seeding method
- Ice nucleation
- Ice nucleation at high pressure
- Ice nucleation from salty water
- Future work

# Classical Nucleation Theory (CNT)

activated process

$$\Delta G(N) = -N|\Delta\mu| + A\gamma$$

$\gamma$  Liquid-solid interfacial  
free-energy at coexistence

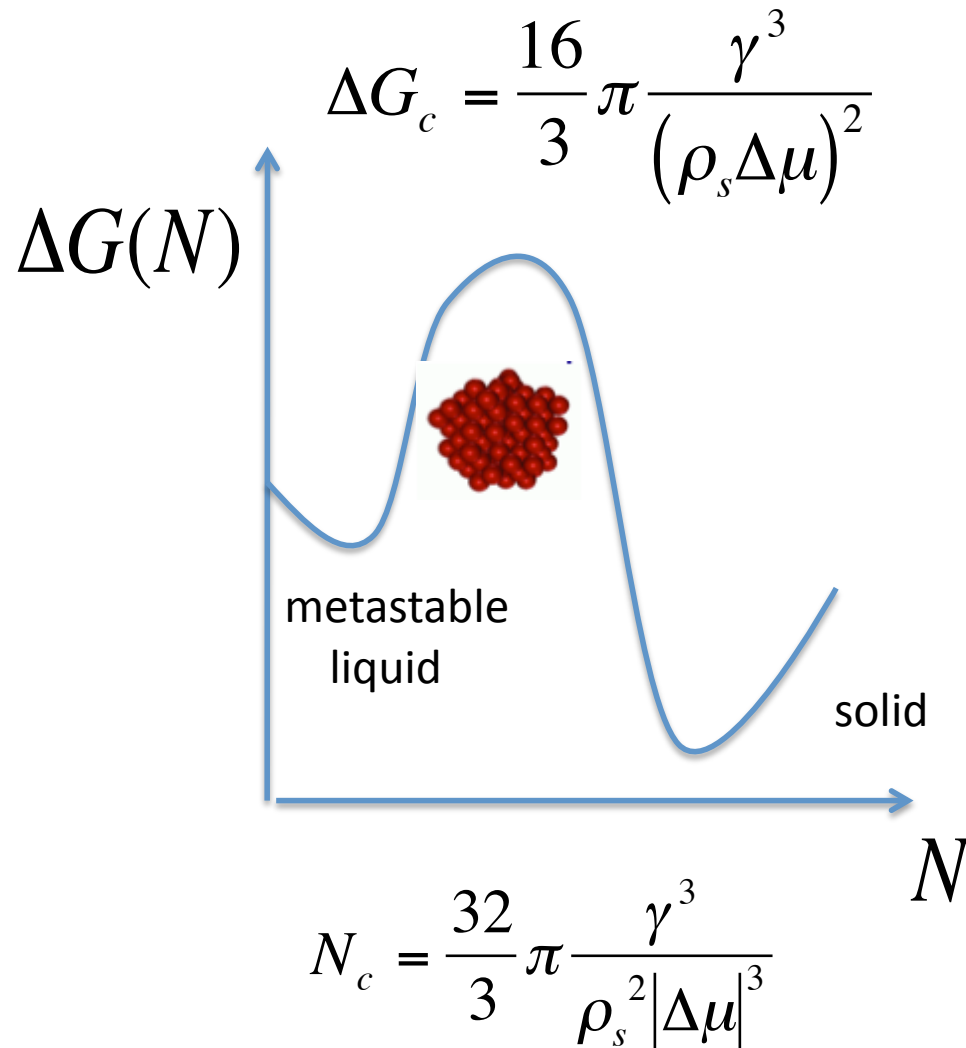


**Order parameter:**  
number of ice molecules  
in the largest crystalline cluster



# Classical Nucleation Theory (CNT)

activated process

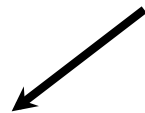


Experimentally it is not possible to observe the critical cluster: too small and short lived

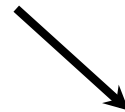
# Steady-state nucleation rate

- J can be measured experimentally
- J can be computed numerically

$$J = A \exp^{-\beta \Delta G_c}$$

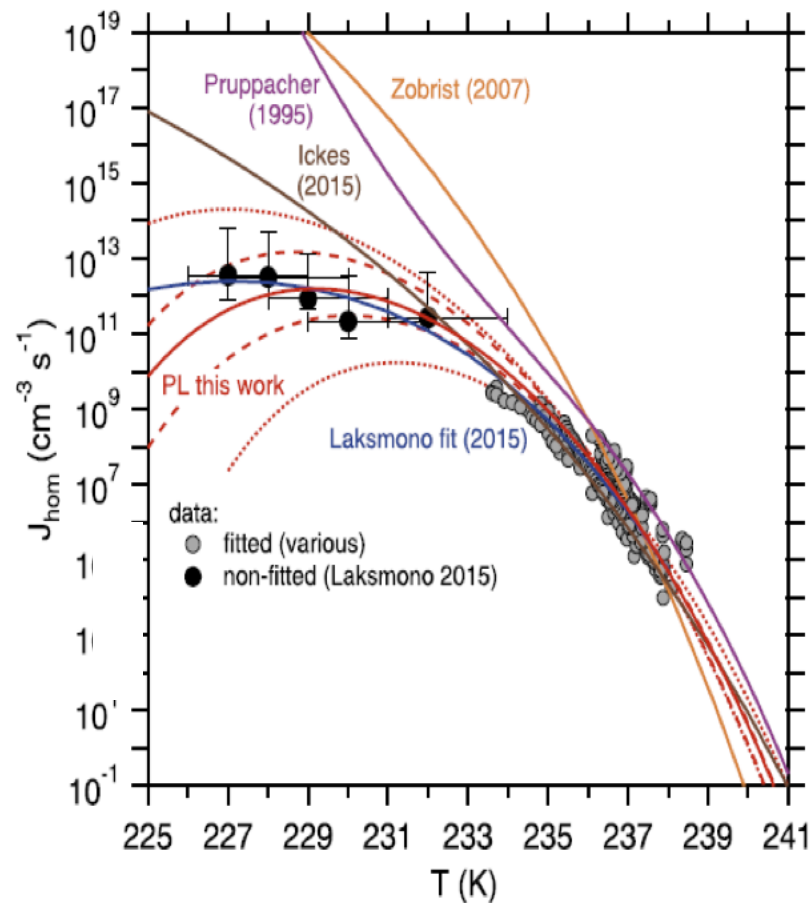


kinetic prefactor  
(flux across the barrier)

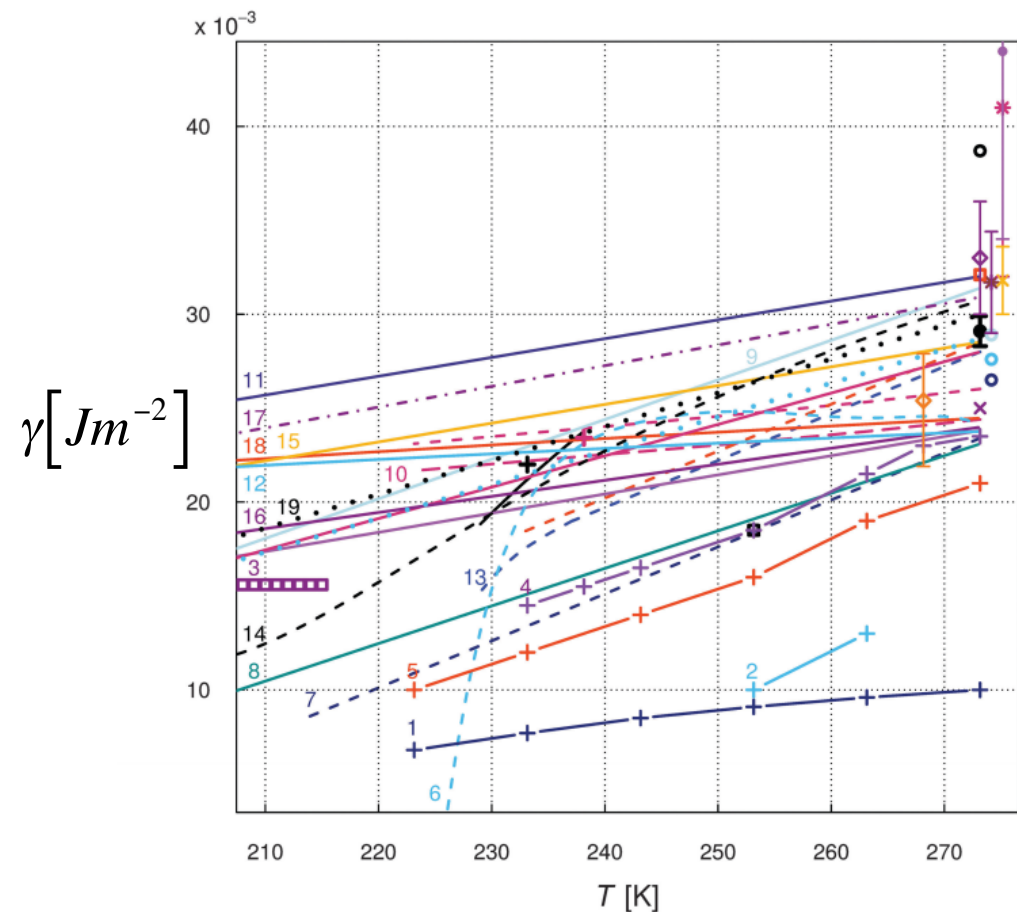


probability to form the critical cluster at  
the top of the free-energy barrier

# Experimental results for $\gamma$ and $J$



Murray, Koop, JCP 2016



Ickes et al PCCP, 2015

Given that an experimental consensus has not yet been reached, computer simulations could play a role

# Seeding method to compute J

- CNT & Numerical simulations
- It is an approximated method, based on the apriori assumption of the cluster structure
- It allows to evaluate nucleation rates within broad thermodynamic conditions

Sanz, Vega, Espinosa, Caballero-Bernal, Abascal, Valeriani JACS (2013)

Espinosa, Vega, Valeriani, Sanz JCP (2016)

# Applying the seeding method

Make use of the CNT expression of the nucleation rate and compute each term by means of computer simulations

$$J = A \exp^{-\beta \Delta G_c}$$

$$J = \sqrt{\frac{|\Delta\mu|}{6\pi\kappa_B T N_c}} \rho_f f^+ \exp\left(-\frac{N_c |\Delta\mu|}{2\kappa_B T}\right)$$



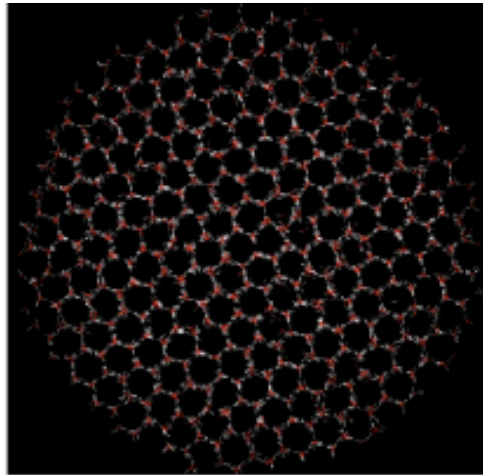
# Applying the seeding method

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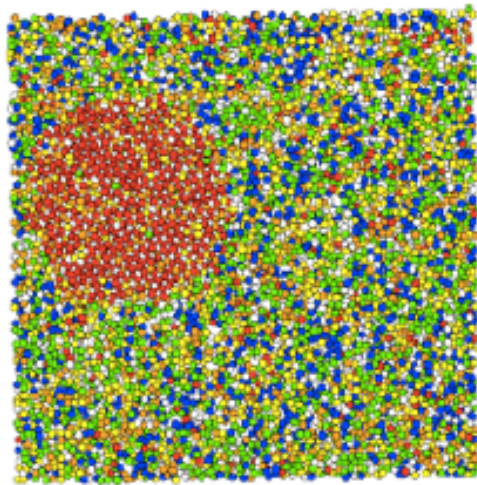
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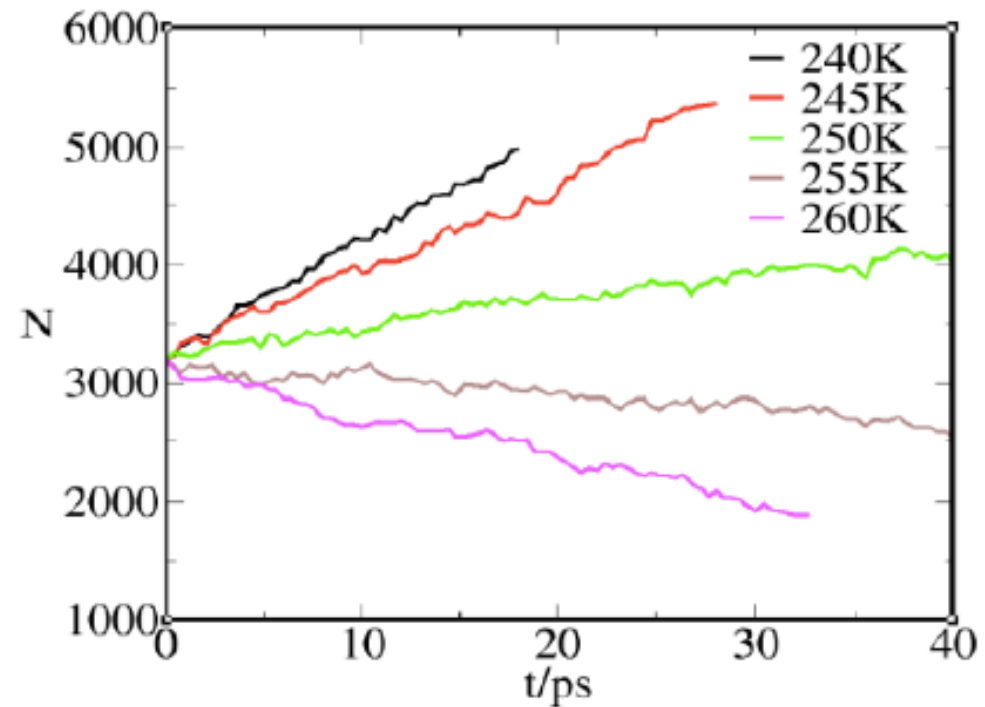
# Computing $N_c$ (critical cluster size)



1- Prepare a crystal seed

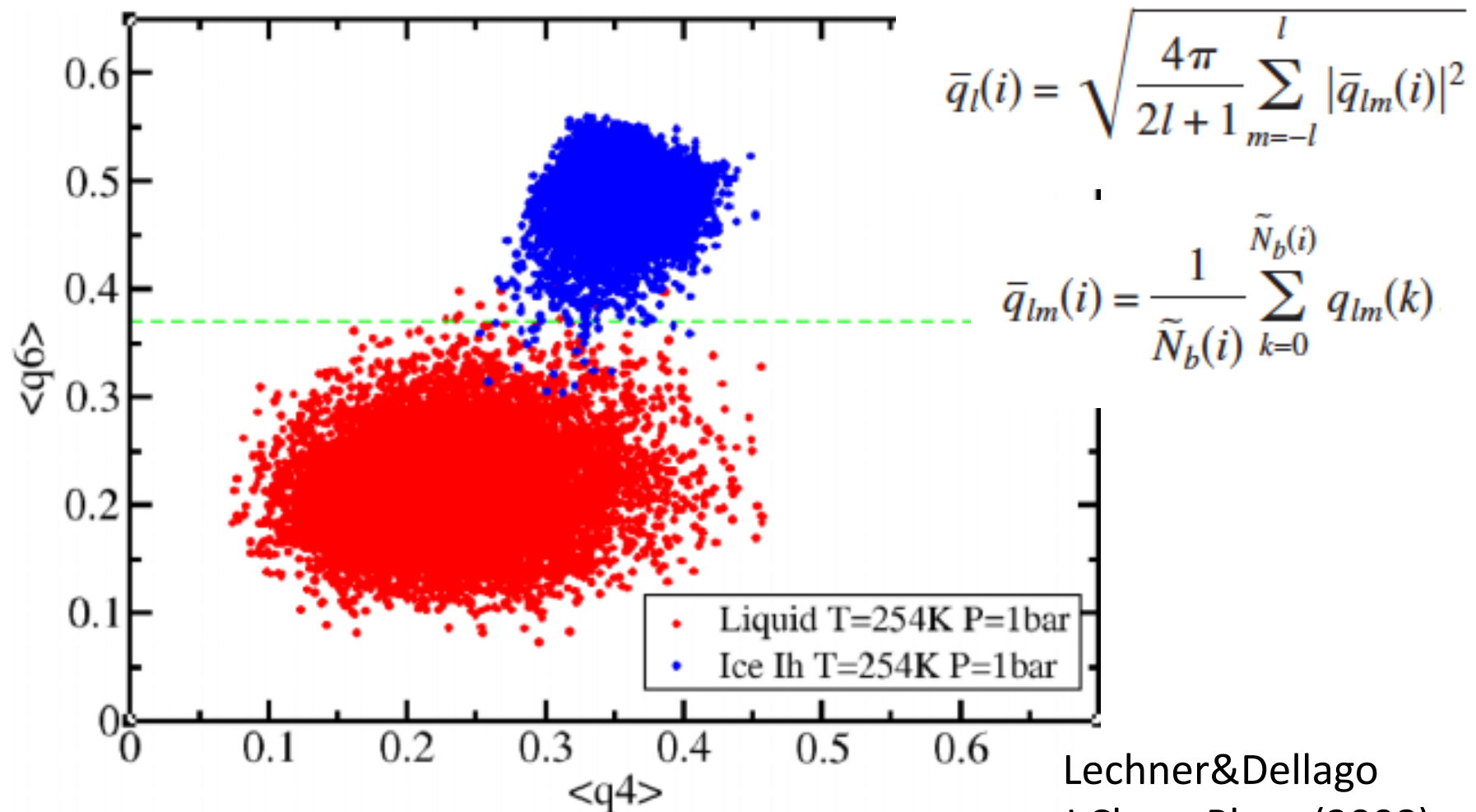


2- Embed the seed in the supercooled melt



3- Launch simulations at different temperatures to find the temperature at which the cluster is critical:  $T_2 = 252.5$  K in this case.

# PARENTHESIS: distinguishing liquid from ice



$$\bar{q}_l(i) = \sqrt{\frac{4\pi}{2l+1} \sum_{m=-l}^l |\bar{q}_{lm}(i)|^2}$$

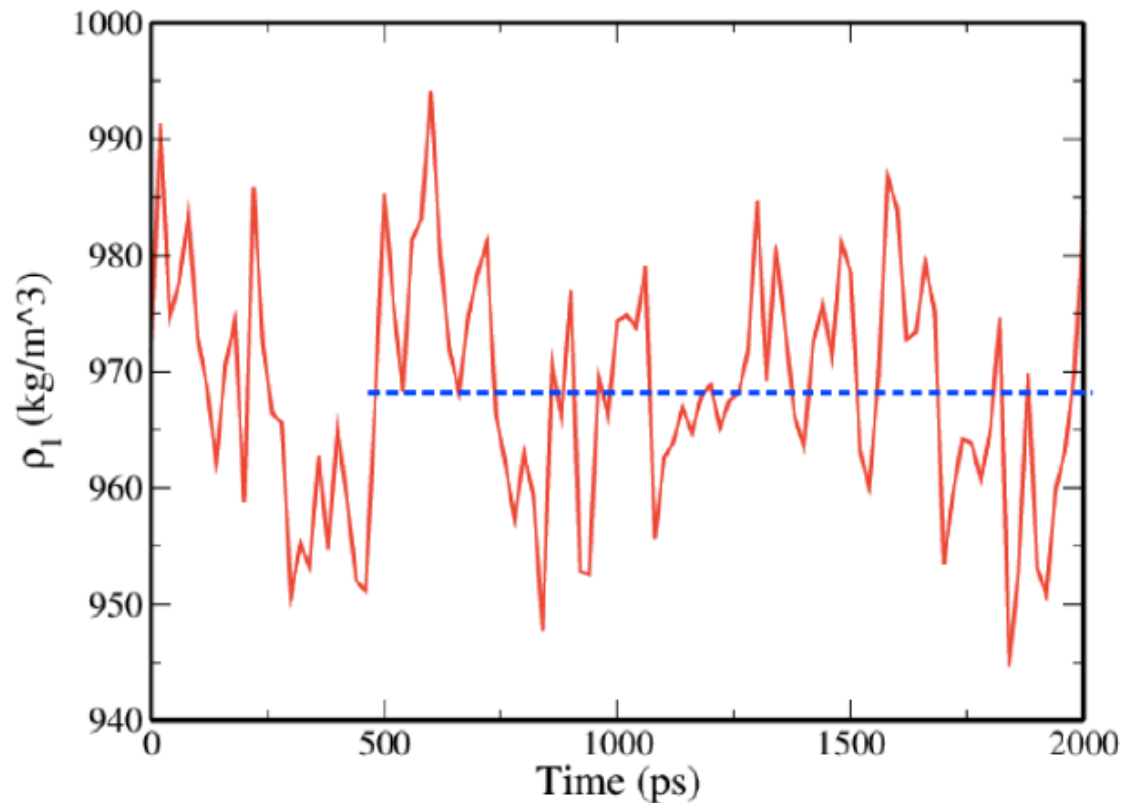
$$\bar{q}_{lm}(i) = \frac{1}{\tilde{N}_b(i)} \sum_{k=0}^{\tilde{N}_b(i)} q_{lm}(k)$$

Lechner&Dellago  
J.Chem.Phys. (2008)

# Computing $\rho_f \Delta\mu f^+$

$$J = \sqrt{\frac{|\Delta\mu|}{6\pi\kappa_B T N_c}} \rho_f f^+ \exp\left(-\frac{N_c |\Delta\mu|}{2\kappa_B T}\right)$$

- $\rho_f$  via a NPT simulations:



# Computing $\rho_f \Delta\mu f^+$

$$J = \sqrt{\frac{|\Delta\mu|}{6\pi\kappa_B T N_c}} \rho_f f^+ \exp\left(-\frac{N_c |\Delta\mu|}{2\kappa_B T}\right)$$

- $\rho_f$  via a NPT simulations:
- $\Delta\mu$  via thermodynamic integrations:



# Thermodynamic integration

$$J = \sqrt{\frac{|\Delta\mu|}{6\pi\kappa_B T N_c}} \rho_f f^+ \exp\left(-\frac{N_c |\Delta\mu|}{2\kappa_B T}\right)$$

By means of thermodynamic integration



$$\Delta\mu \left\{ \begin{array}{l} - \frac{\mu_{ice}(T_2, p)}{k_B T_2} = \cancel{\frac{\mu_{ice}(T_1, p)}{k_B T_1}} - \int_{T_1}^{T_2} \frac{H_{ice}(T)}{Nk_B T^2} dT \\ \frac{\mu_{water}(T_2, p)}{k_B T_2} = \cancel{\frac{\mu_{water}(T_1, p)}{k_B T_1}} - \int_{T_1}^{T_2} \frac{H_{water}(T)}{Nk_B T^2} dT \end{array} \right.$$

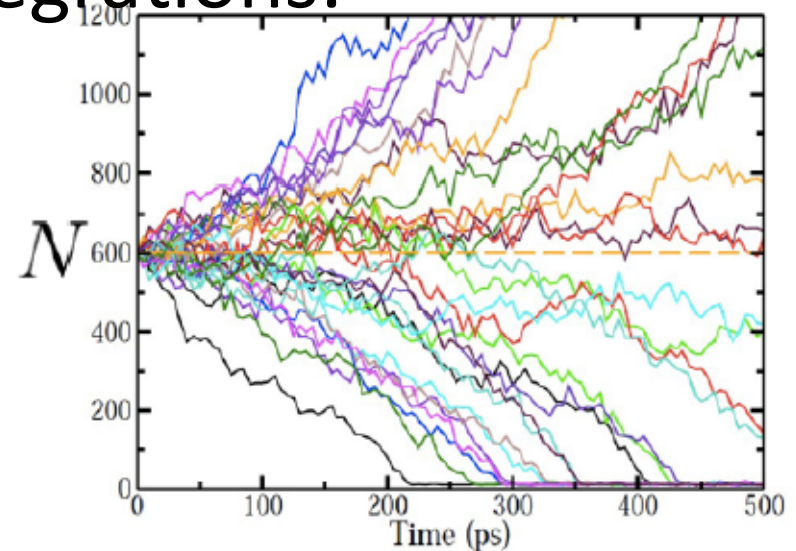
# Computing $\rho_f \Delta\mu f^+$

$$J = \sqrt{\frac{|\Delta\mu|}{6\pi\kappa_B T N_c}} \rho_f f^+ \exp\left(-\frac{N_c |\Delta\mu|}{2\kappa_B T}\right)$$

- $\rho_f$  via a NPT simulations:

- $\Delta\mu$  via thermodynamic integrations:

- $f^+ \equiv \left\langle \frac{(N - N_c)^2}{2t} \right\rangle$  Auer&Frenkel  
Nature (2001)



Now we compute the rate  $J$

# Why do we need a supercomputer?

To compute the nucleation rate  $J$  for three crystalline clusters (i.e. at three super-coolings), we have used 300000 CPU hours.

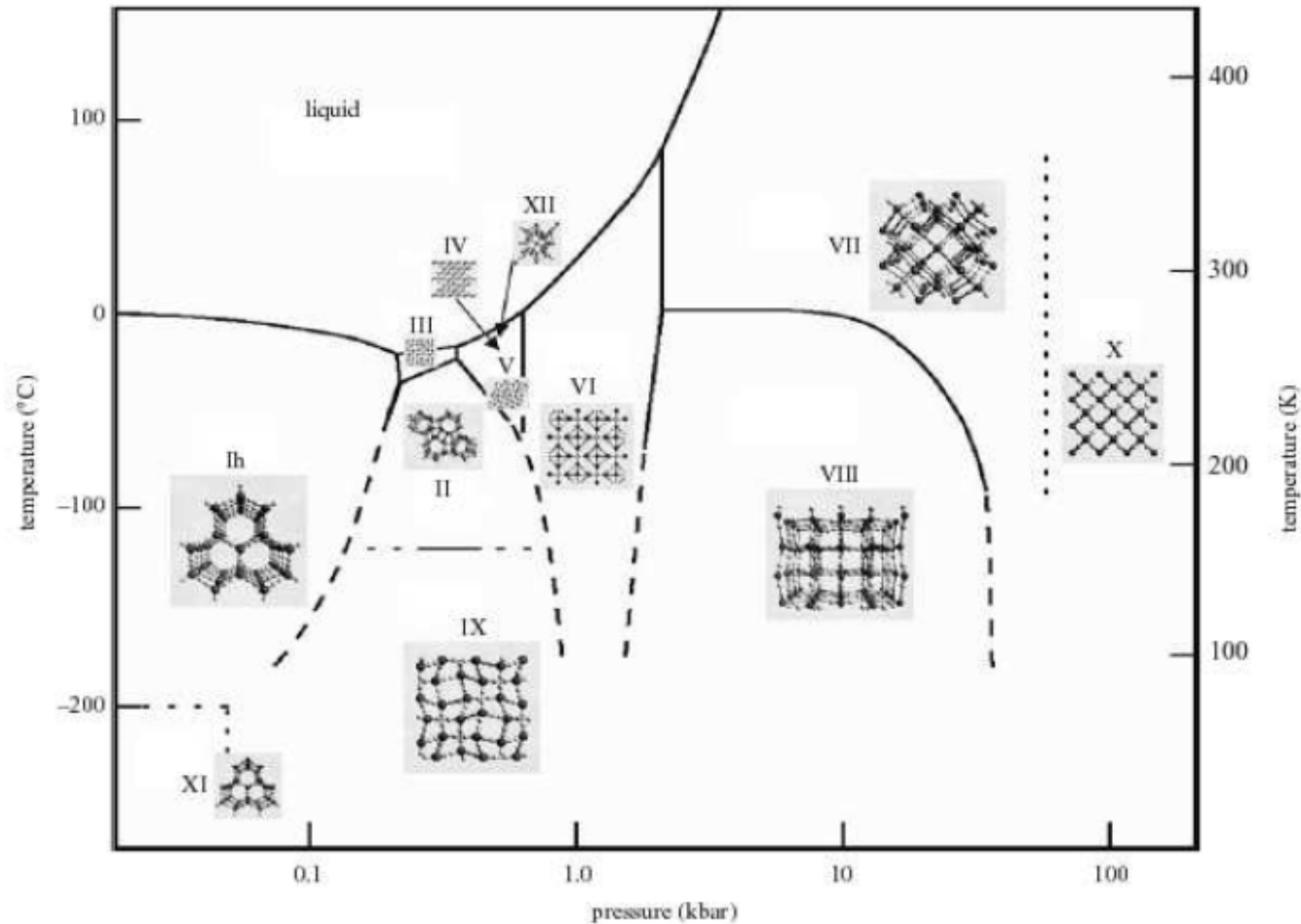
Running only on 1 processor, would take about 30 years!

Supercomputers are needed to speed up our work

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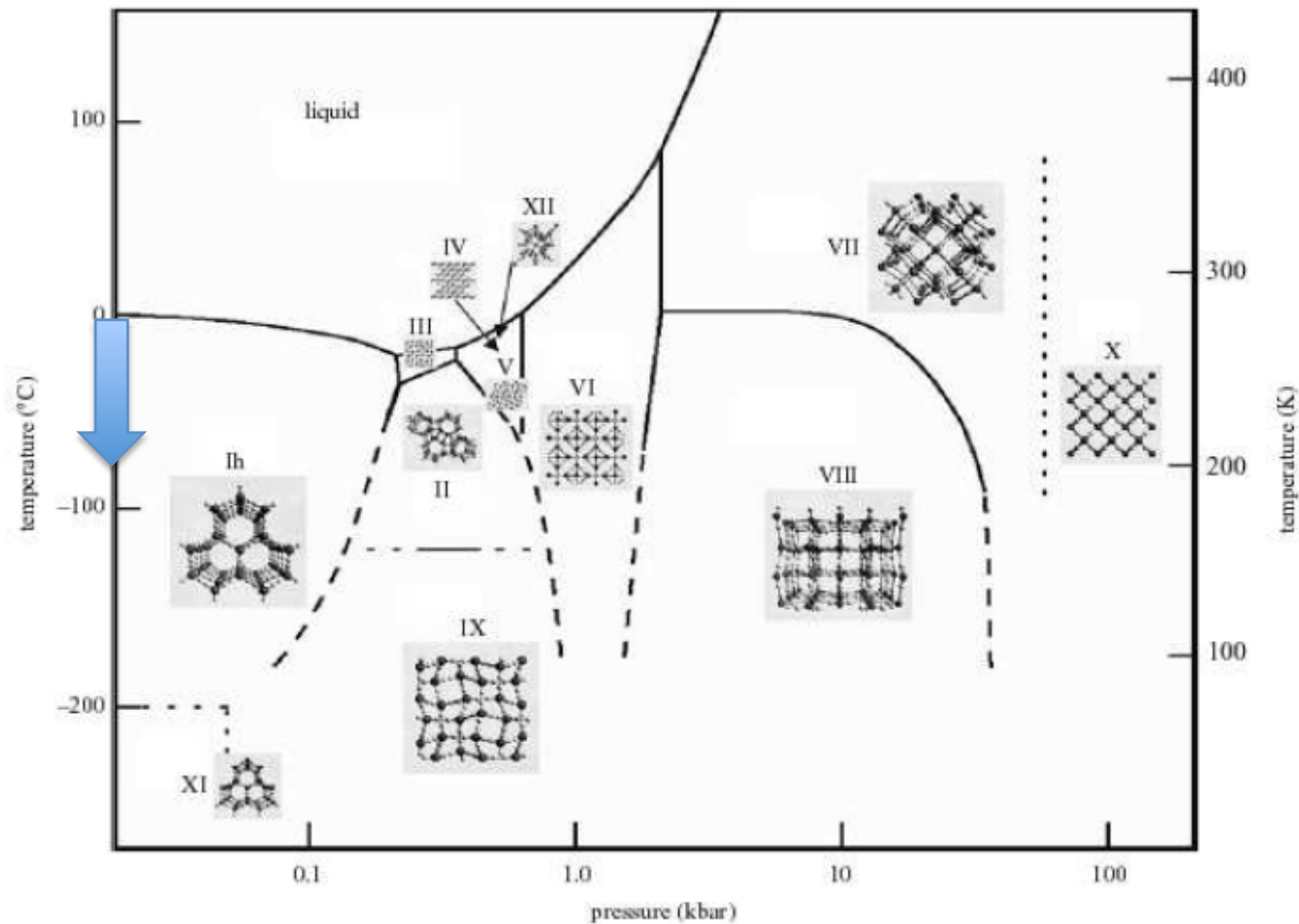
# Water's phase diagram



J.Finney, Phyl.Trans.R.Soc.Lond.B,(2004)

1900 Tammann , 1912 Bridgman, 1968 Whalley , 2009 Finney et al.

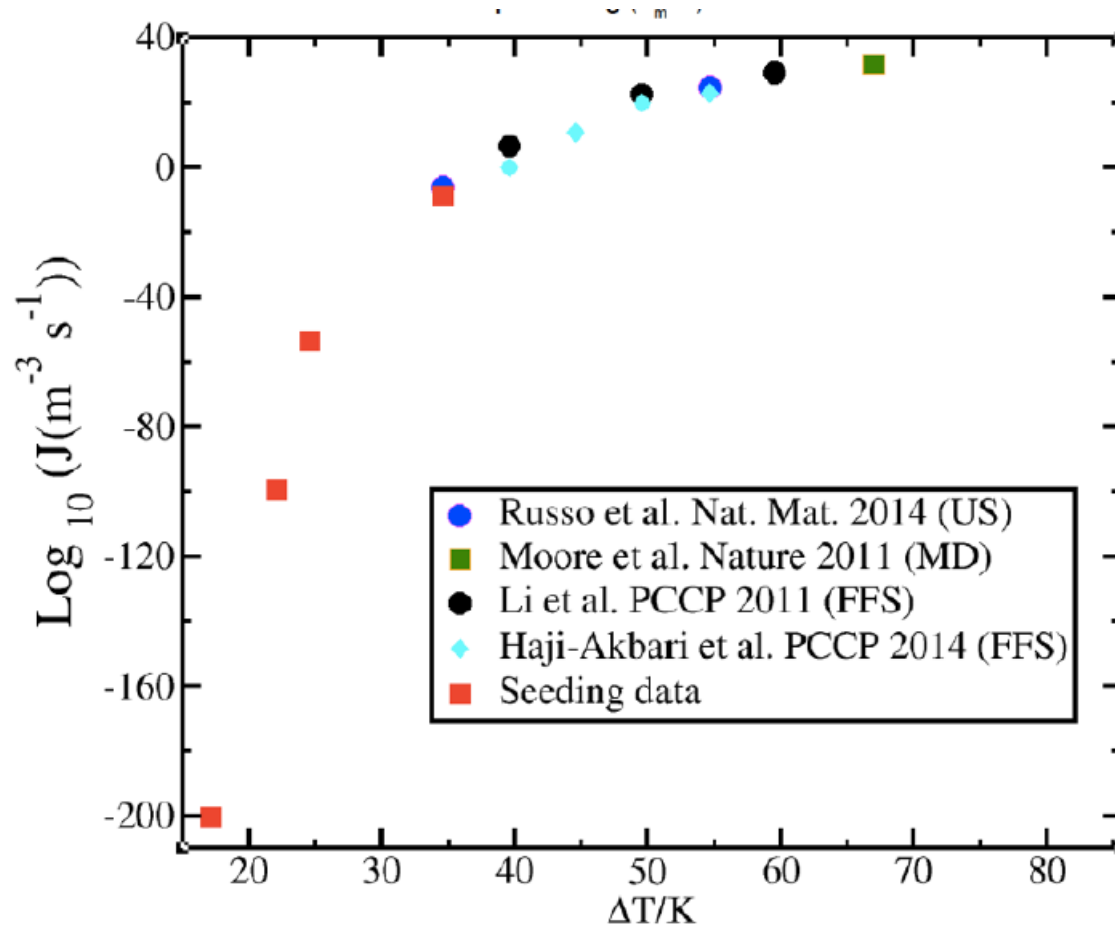
# Water's phase diagram



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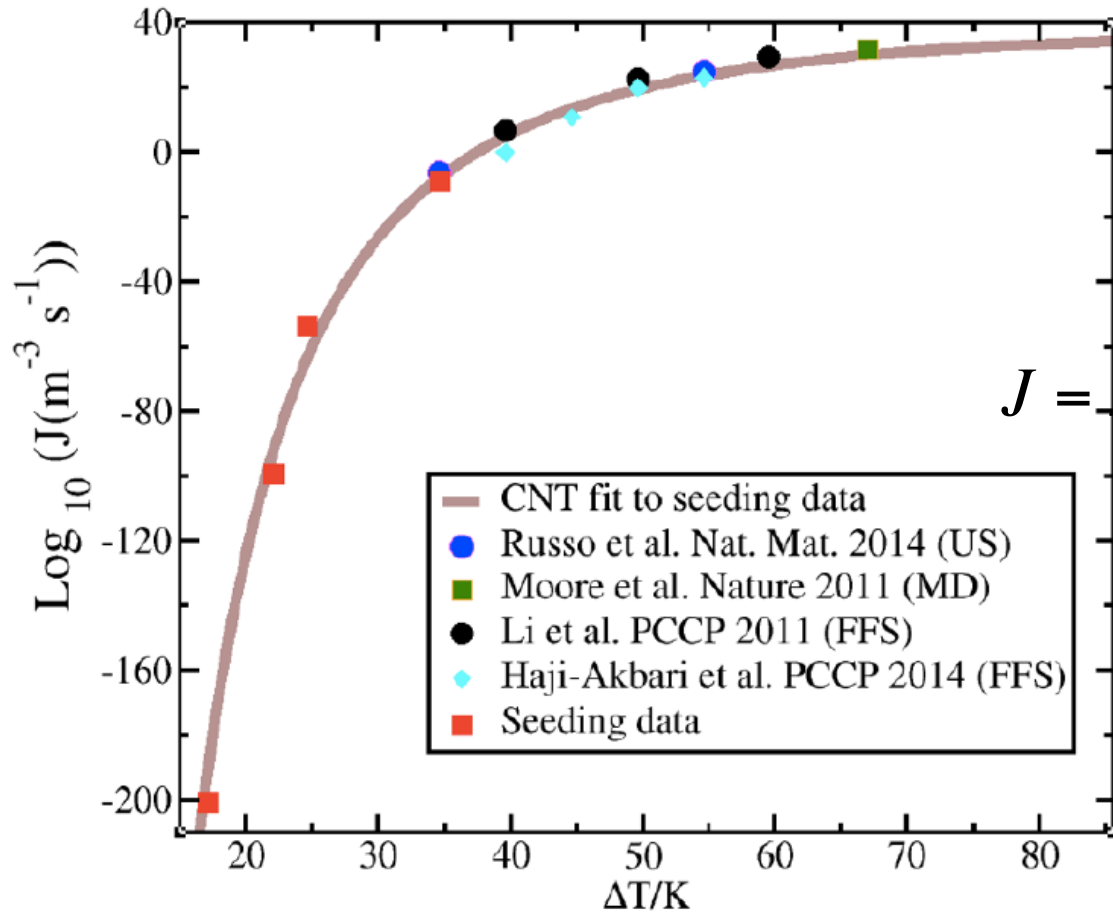
# Ice nucleation rate at 1 bar



mW water

Molinero&Moore,  
J.Phys.Chem.B(2009)

# Ice nucleation rate at 1 bar



Interpolating all quantities within a broad T range

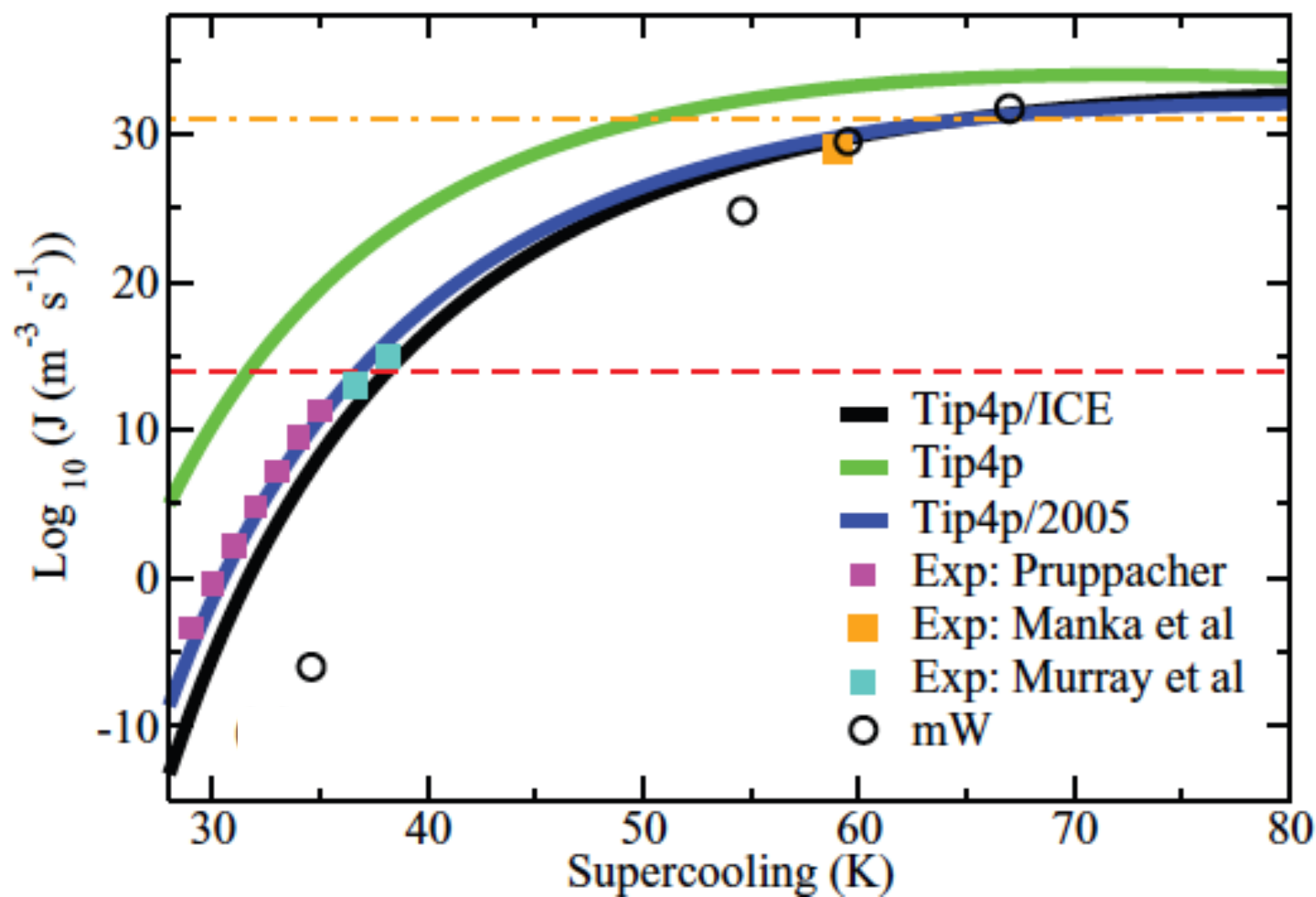
$$J = \sqrt{\frac{|\Delta\mu|}{6\pi\kappa_B T N_c}} \rho_s f^+ \exp\left(-\frac{N_c |\Delta\mu|}{2\kappa_B T}\right)$$

according to CNT

$$f^+ \equiv \frac{24 D N_c^{2/3}}{\lambda^2}$$



# Comparing several water models to experiments

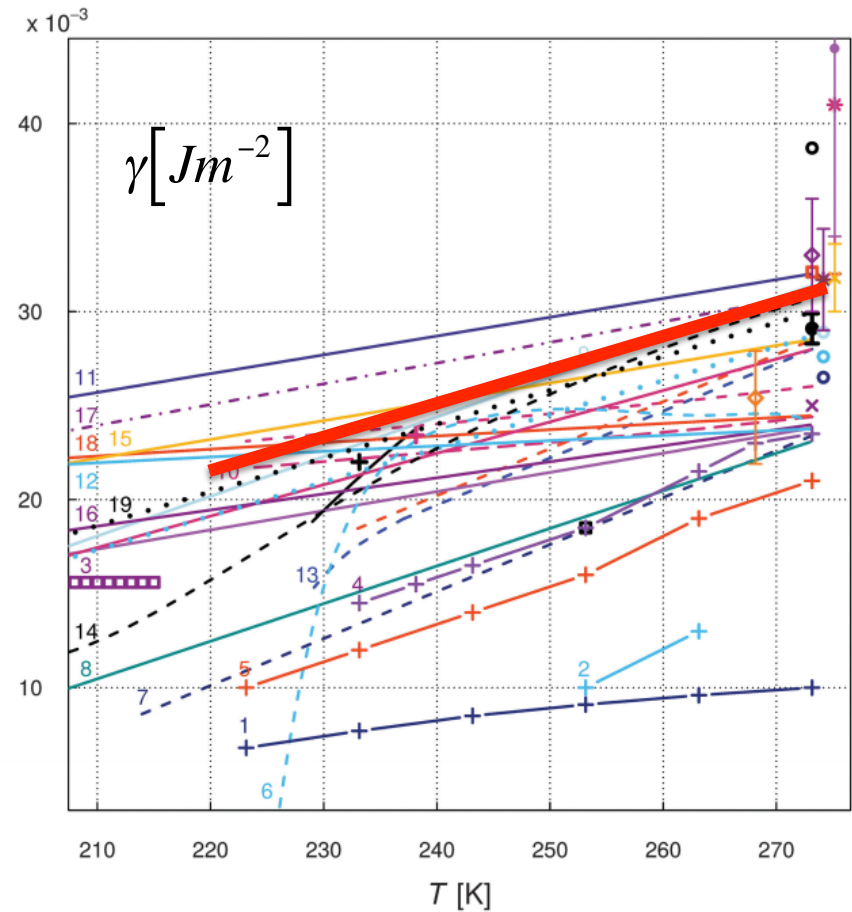
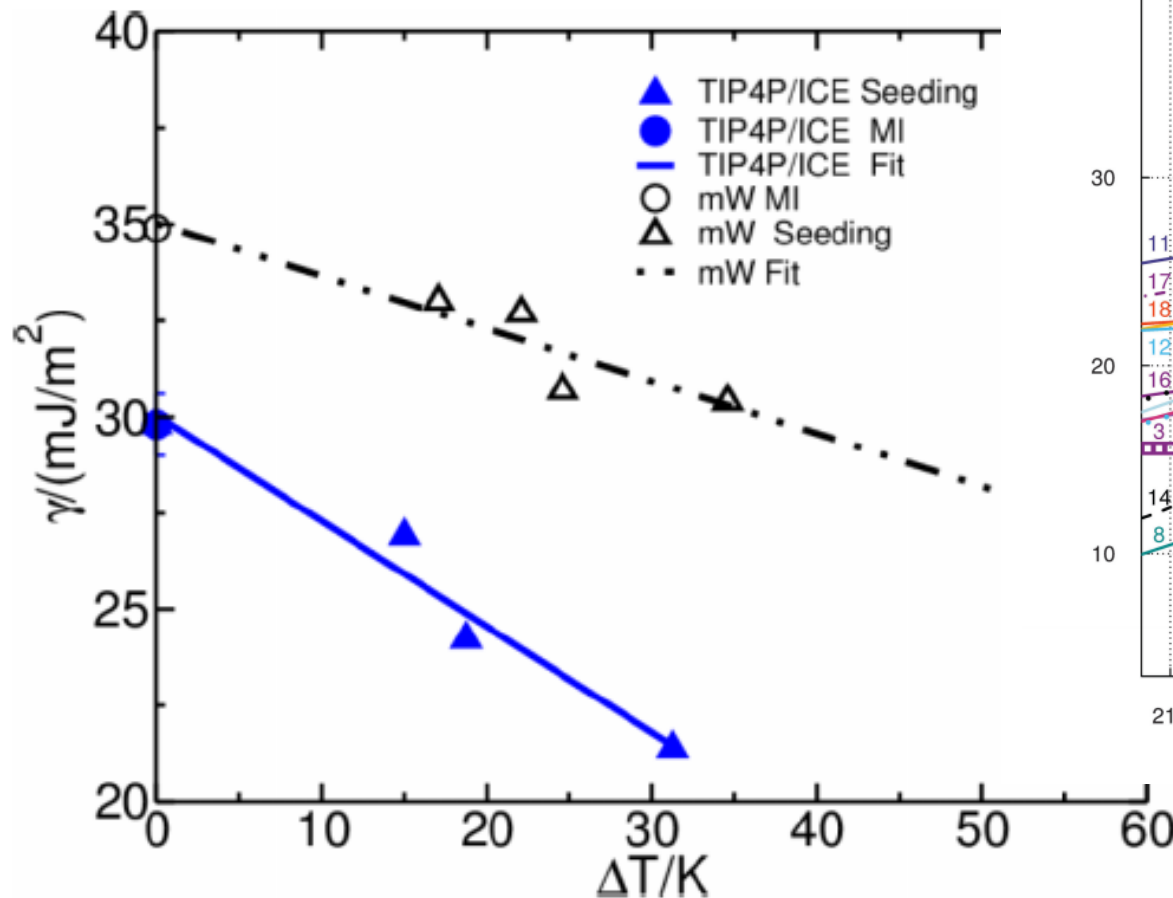


E.Sanz, C.Vega, J.Espinosa, R.Bernal, J.Abascal and C.Valeriani JACS (2013)

J.Espinosa, E.Sanz, C.Valeriani and C.Vega JCP (2014)

# Consequence 1: liquid-ice interfacial free-energy

$$N_c = \frac{32}{3} \pi \frac{\gamma^3}{\rho_s^2 |\Delta\mu|^3}$$



Ickes et al PCCP, 2015

Espinosa, Navarro, Sanz, Valeriani, Vega JCP (2016)

# Consequence 2:

## Ice Ih versus ice Ic

### Extent and relevance of stacking disorder in "ice I<sub>c</sub>"

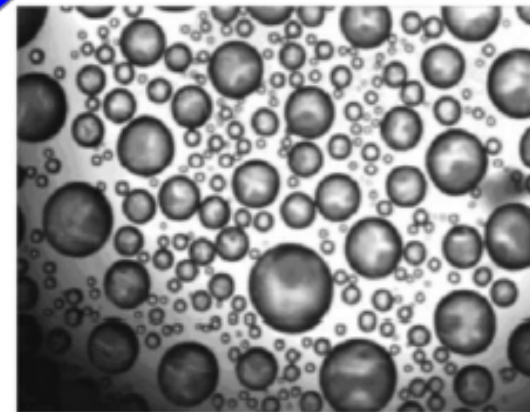
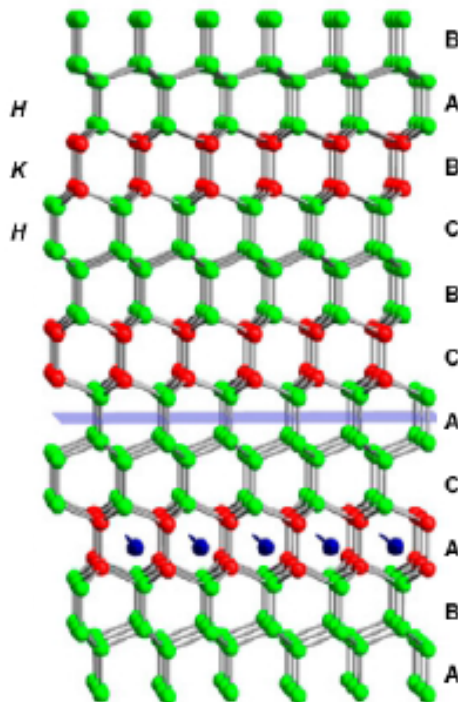
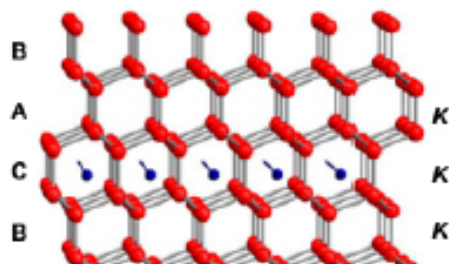
Werner F. Kuhs<sup>a,1</sup>, Christian Sippel<sup>a,b</sup>, Andrzej Falenty<sup>a</sup>, and Thomas C. Hansen<sup>b</sup>

(2012)

### Structure of ice crystallized from supercooled water

Tamsin L. Malkin<sup>a</sup>, Benjamin J. Murray<sup>a,1</sup>, Andrey V. Brukhno<sup>b</sup>, Jamshed Anwar<sup>c</sup>, and Christoph G. Salzmann<sup>d,2</sup>

(2012)

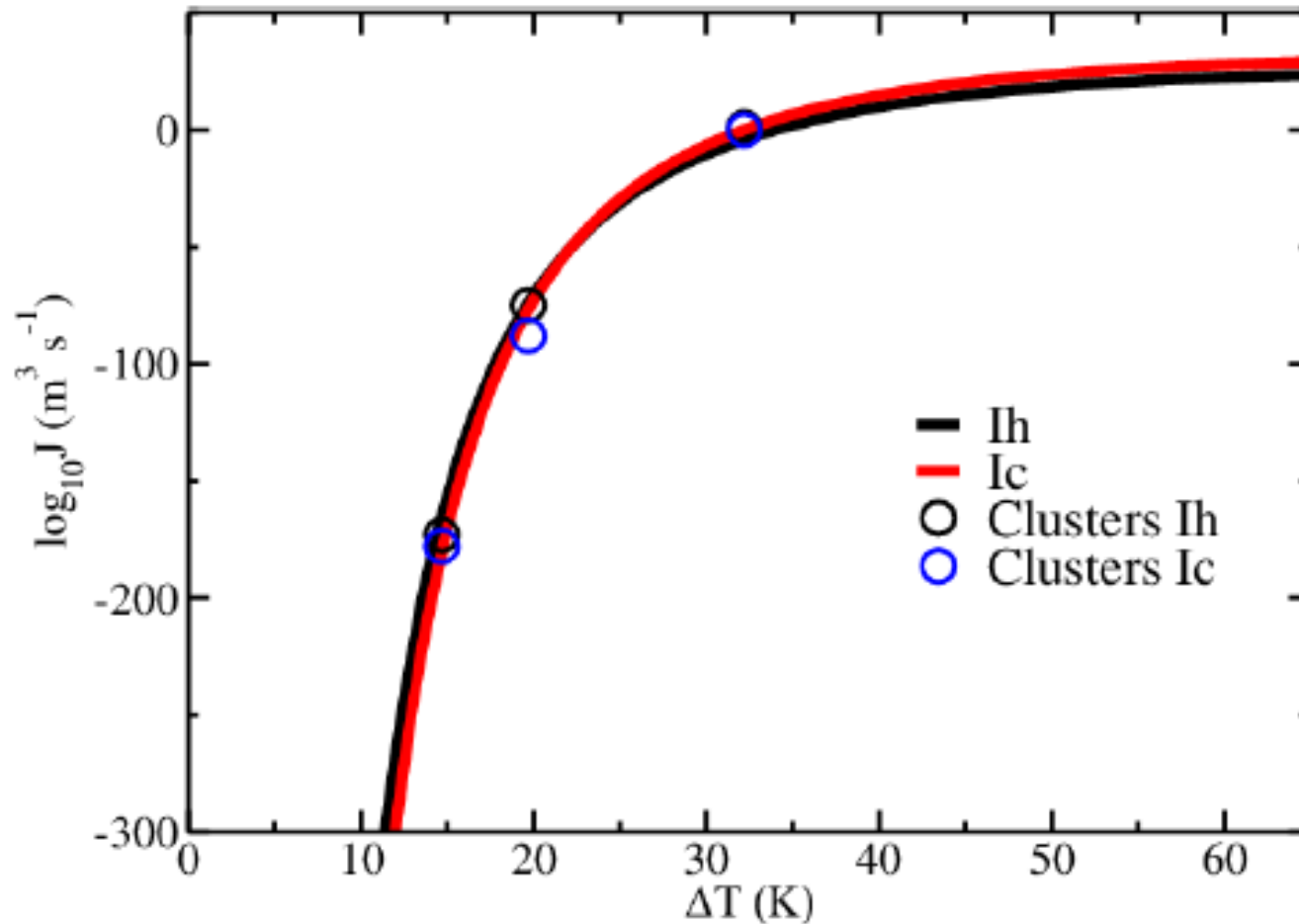


As observed in  
computer simulations  
with mW water

Hudait, Qiu, Lupi and Molinero  
PCCP(2016)

B.J.Murray et al. "*Kinetics of the homogeneous freezing of water*"  
Physical Chemistry Chemical  
Physics. (2010)

# Ice nucleation rate at 1 bar



TIP4P/ICE water

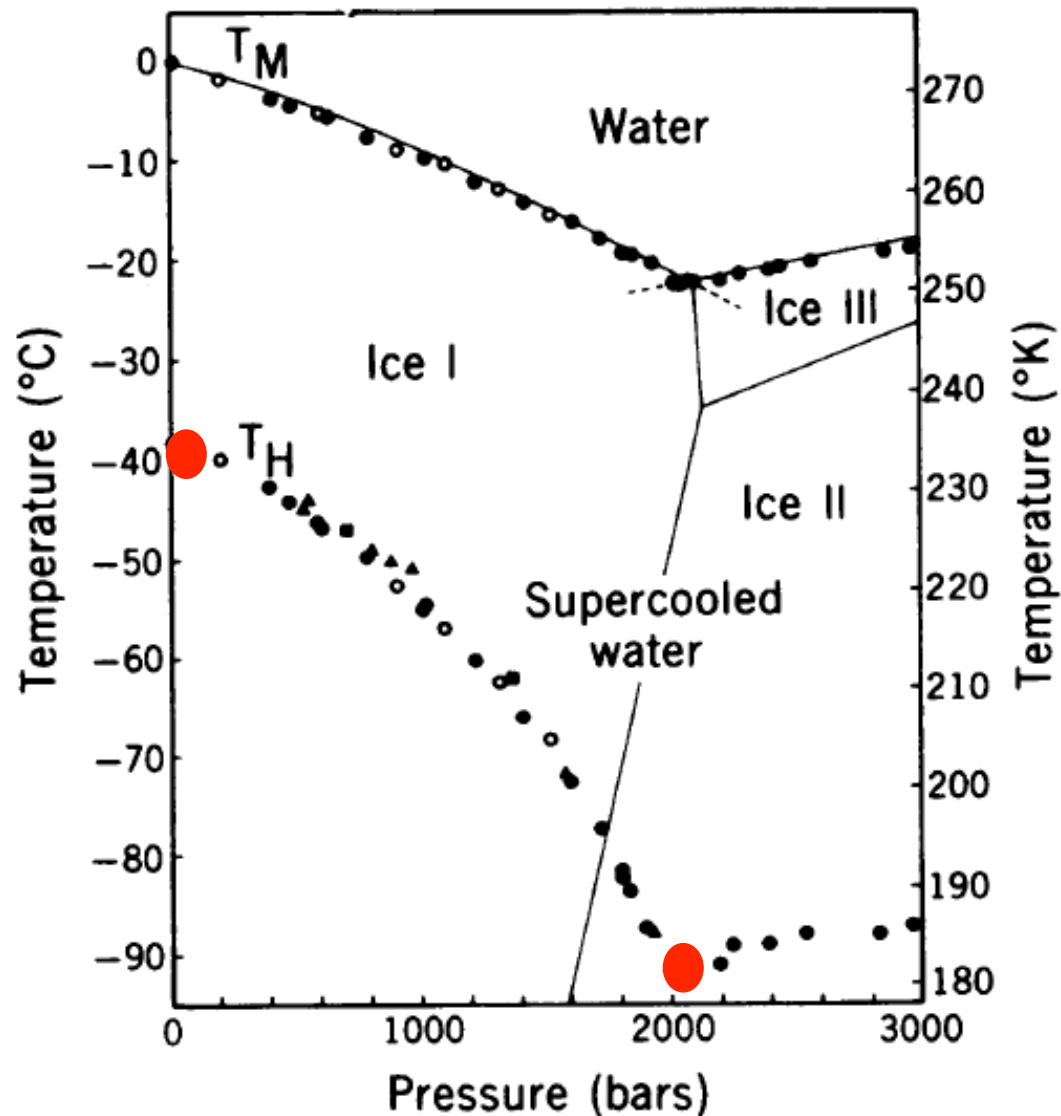
Ice Ih/Ic stacking faults during ice nucleation

Zaragoza, Conde, Espinosa, Valeriani, Vega, Sanz JCP (2015)

# Outline

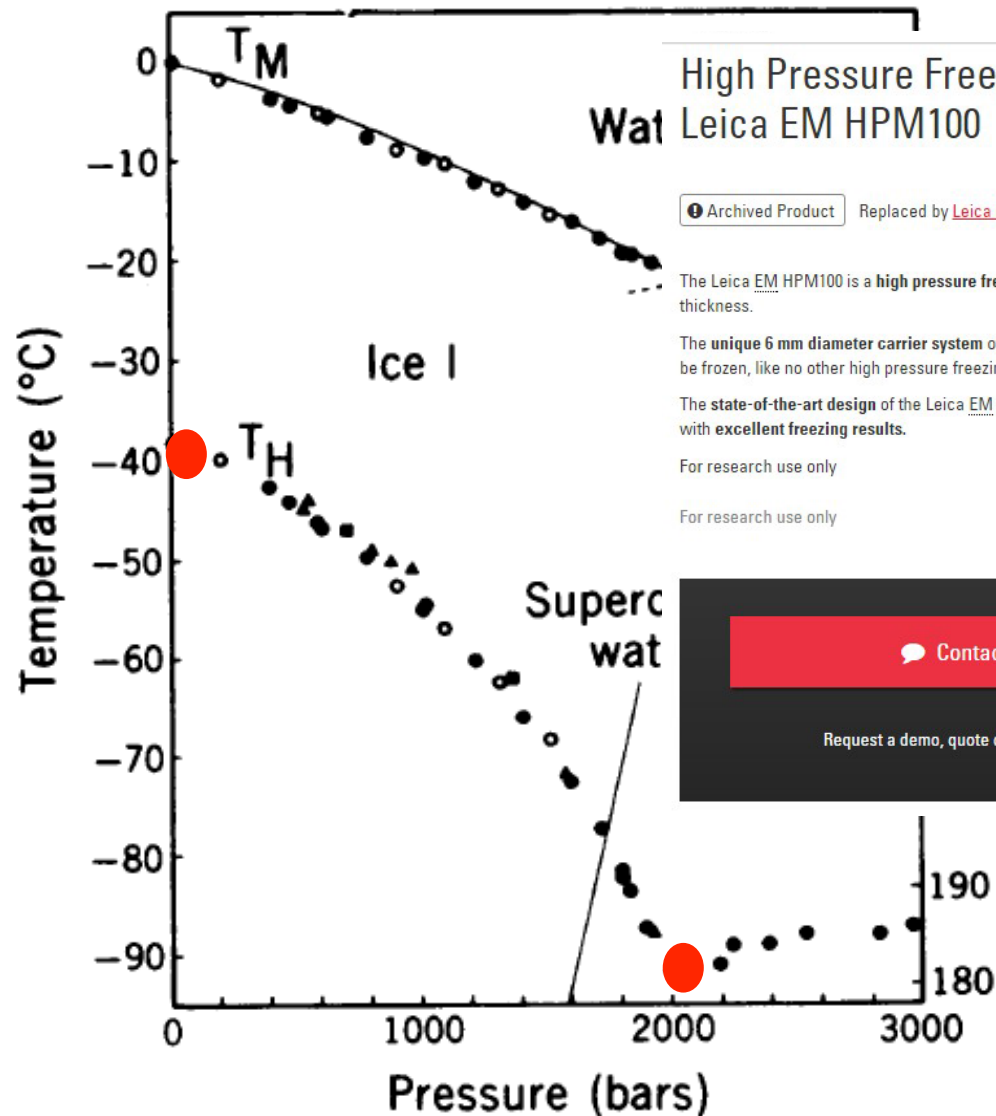
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# Effect of pressure on ice nucleation



Kano, Speedy, Angell Science (1975)

# Effect of pressure on ice nucleation



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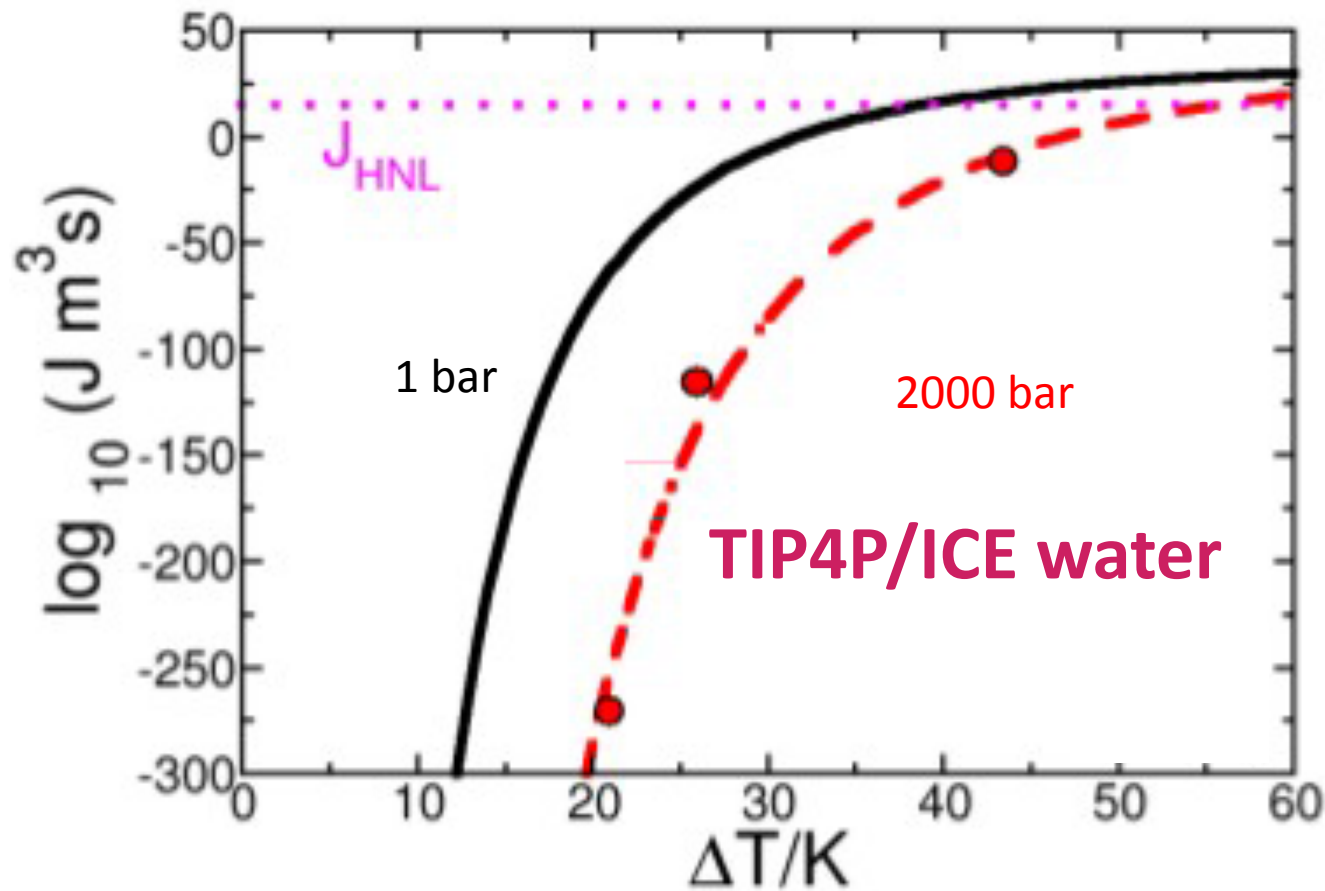
Request a demo, quote or ask a question



## Why does it work?

Kano, Speedy, Angell Science (1975)

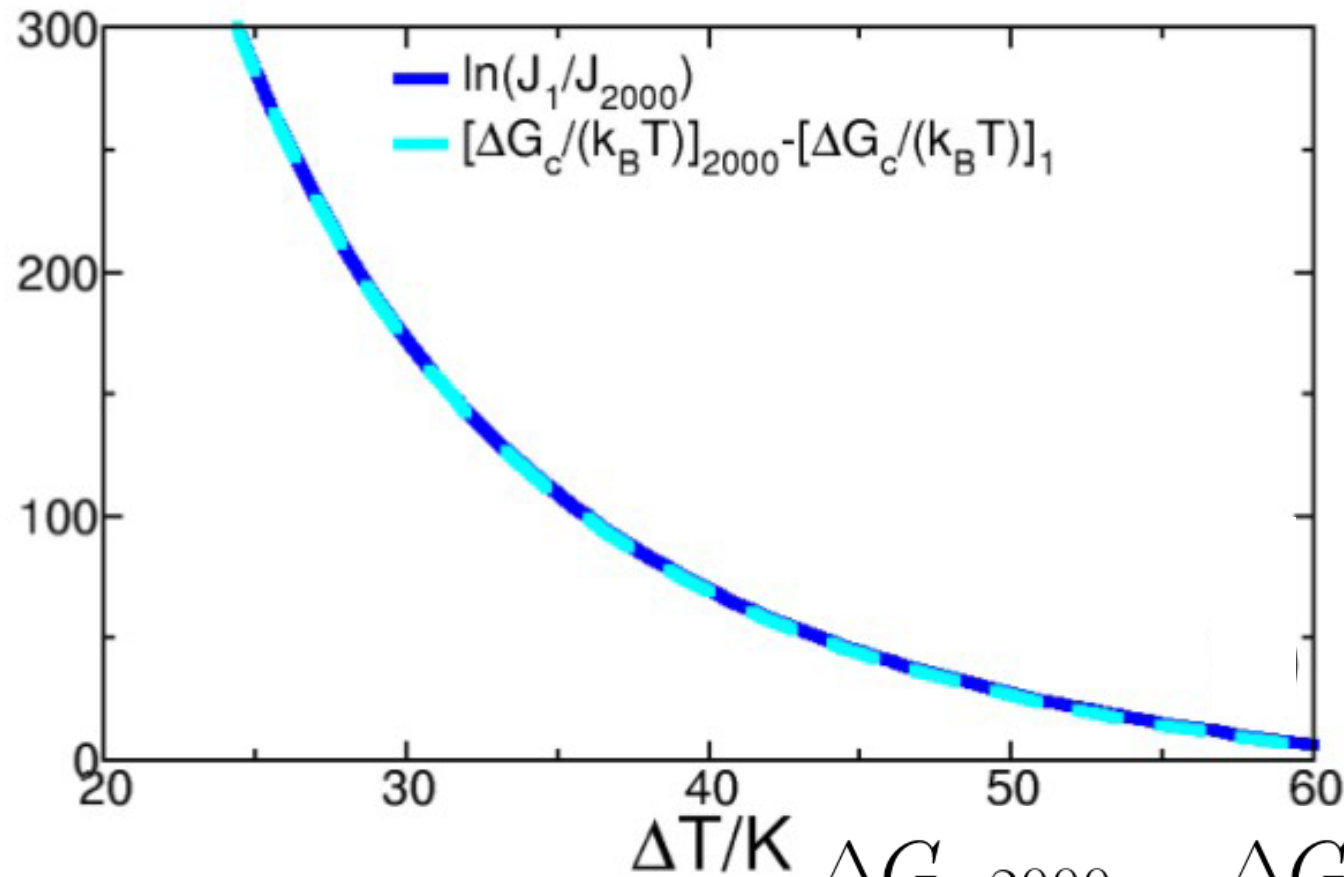
# Homogeneous nucleation rate of ice Ih from supercooled water





# What does it depend on?

$$J = Ae^{-\Delta G_c / K_B T}$$



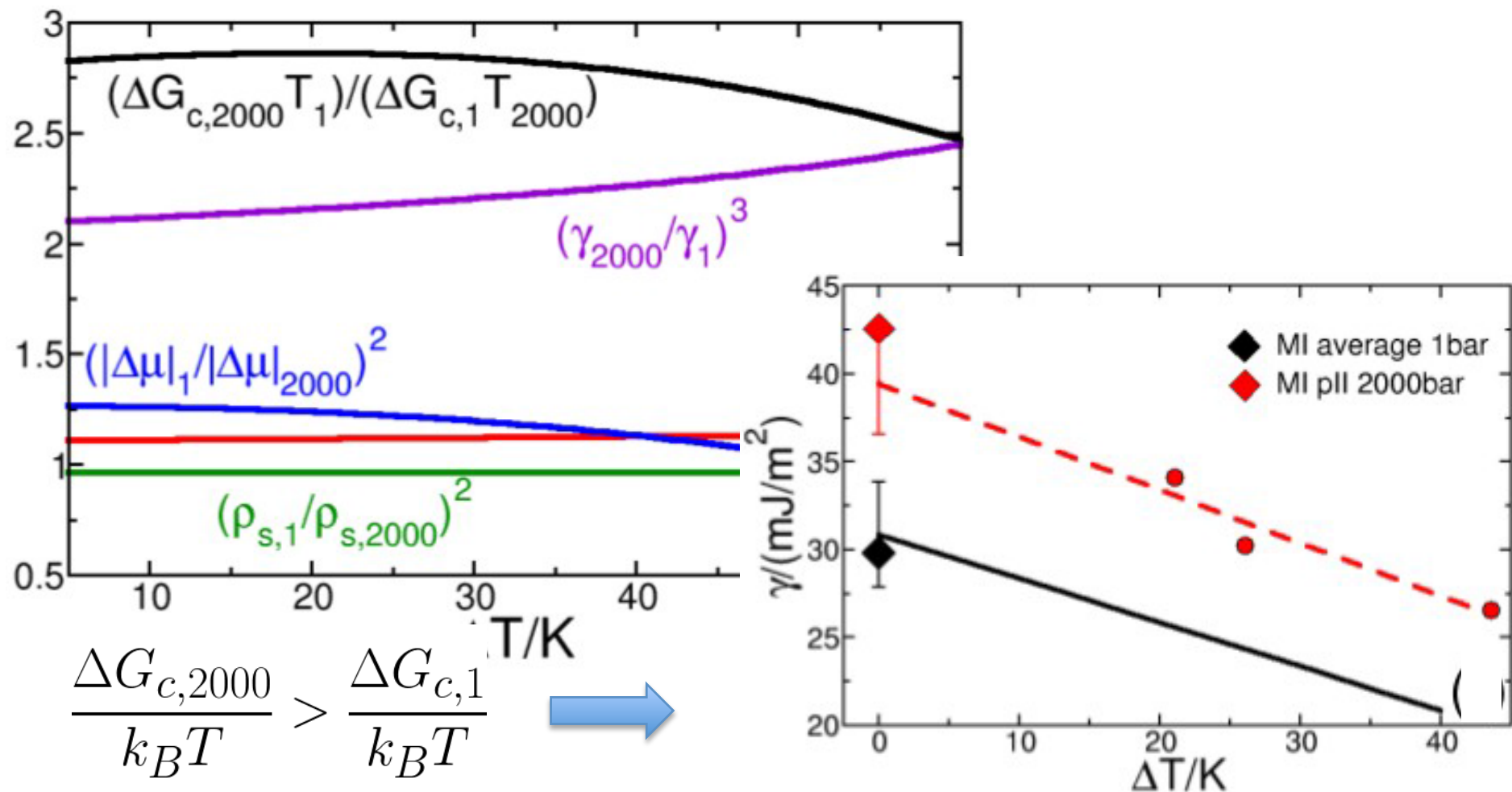
$$A_{2000} \approx A_1$$



$$\frac{\Delta G_{c,2000}}{k_B T} > \frac{\Delta G_{c,1}}{k_B T}$$

# What does it depend on?

$$\frac{\Delta G_c}{k_B T} = \frac{16\pi\gamma^3}{3\rho_s^2|\Delta\mu|^2 k_B T}$$



# Conclusions

High pressure could be a way to supercool water avoiding ice nucleation

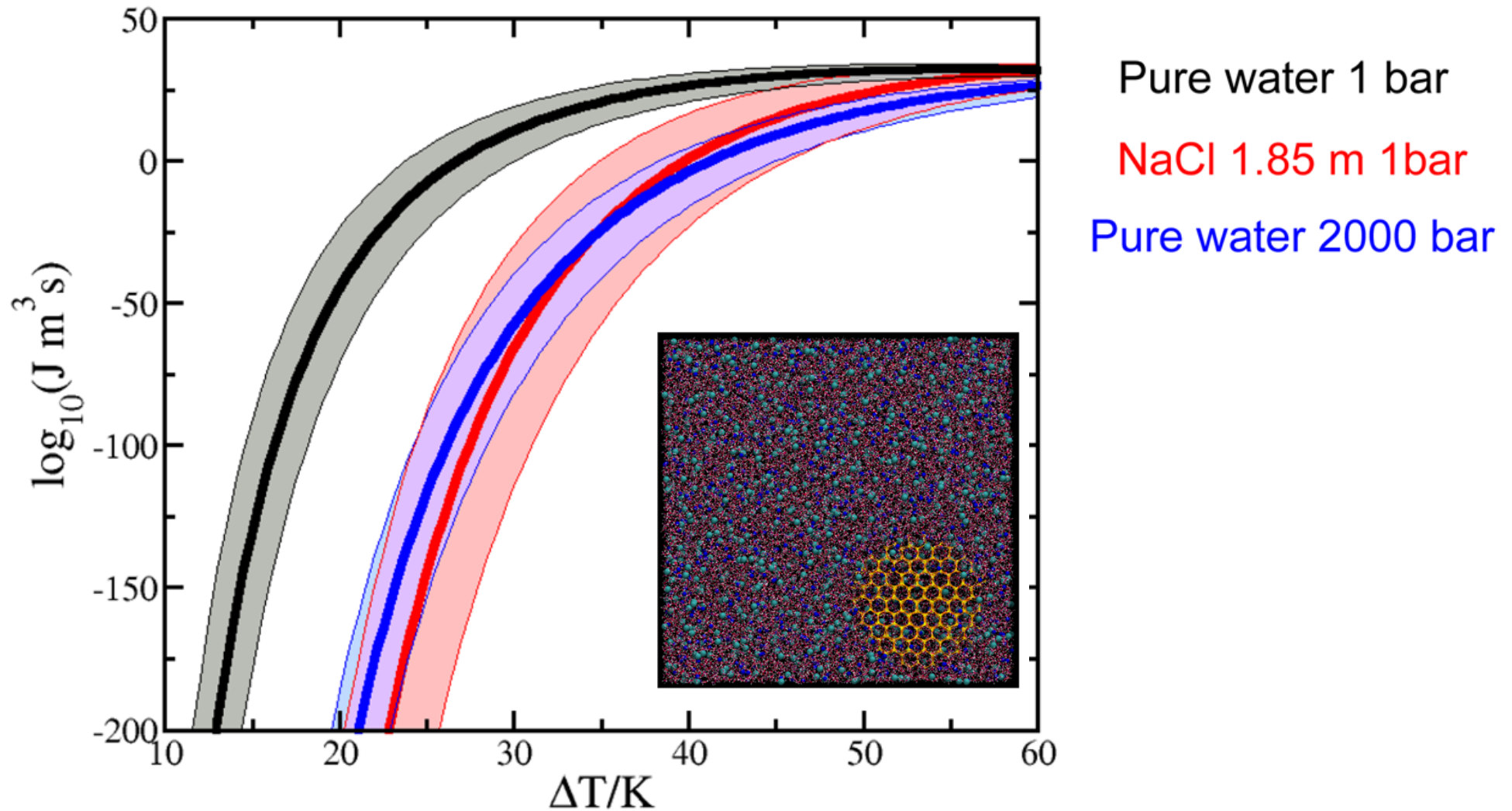
Given that the liquid-solid interfacial free-energy increases with pressure

Espinosa, Zaragoza, Rosales, Navarro, Valeriani, Vega and Sanz, PRL (2016)

# Outline

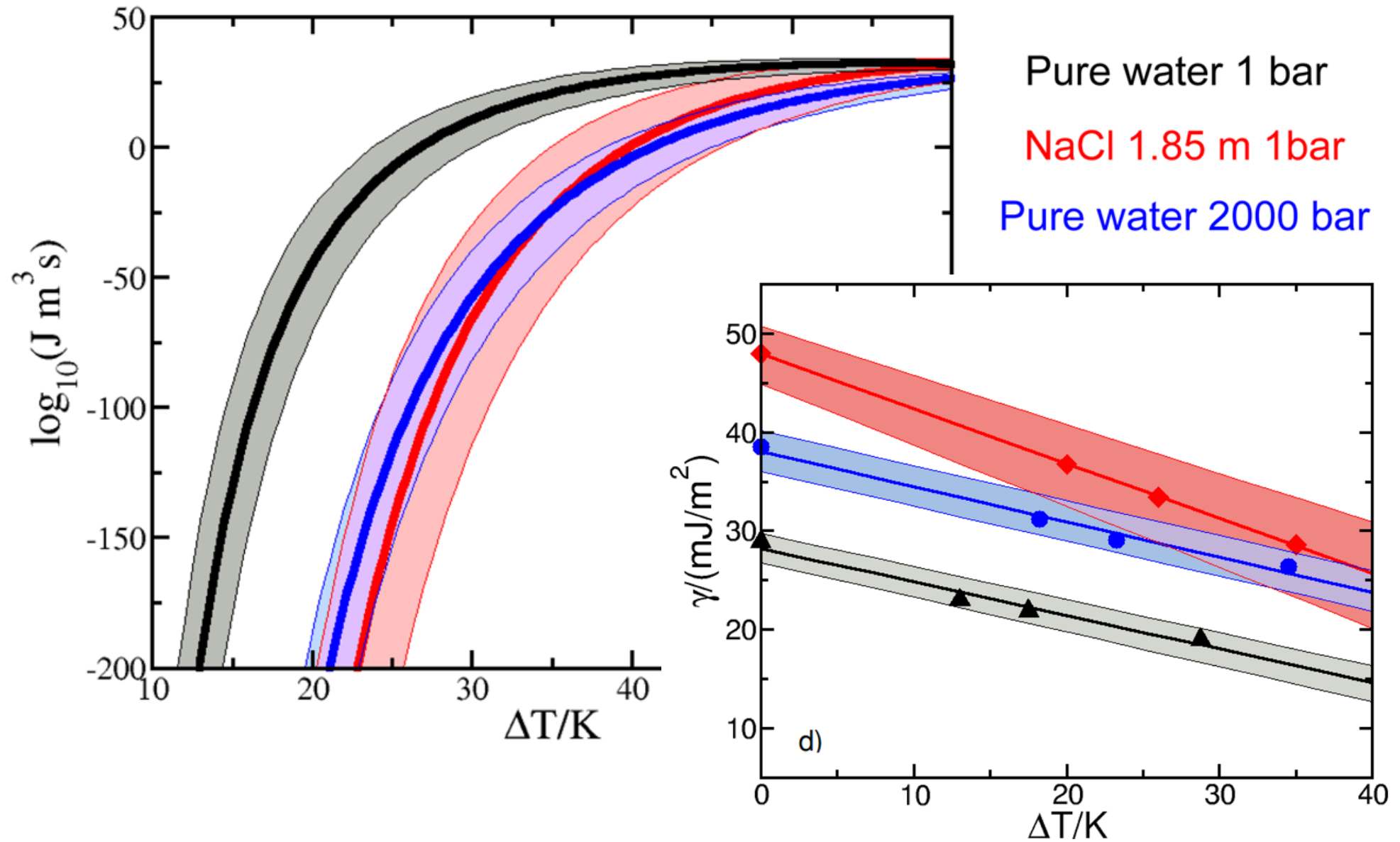
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# Ice nucleation from salty water



Espinosa, Soria, Ramirez, Valeriani, Vega and Sanz, J.Phys.Chem.Lett.(2017)

# Ice nucleation from salty water



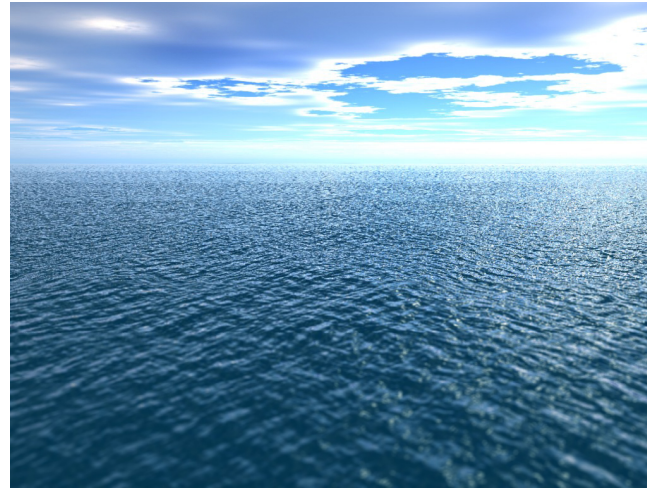
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# Future work:

At -20K

considering all water  
in the hydrosphere



the average time to form homogeneously  
a nucleus would be  $10^{66}$  years!

Maybe bulk ice nucleation in Nature is not homogeneous...

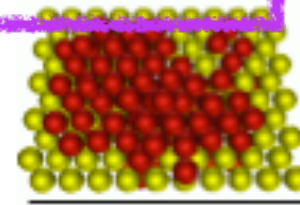
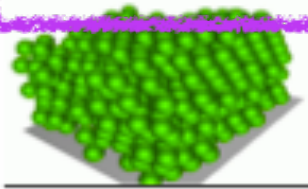
E.Sanz, C.Vega, J.Espinosa, R.Bernal, J.Abascal and C.Valeriani **JACS** (2013)



# Homogeneous versus heterogeneous nucleation



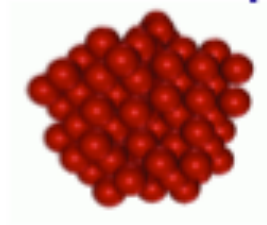
heterogeneous nucleation



crystal seeded by a smooth hard wall

crystal seeded by a regular template

homogeneous nucleation



spontaneous AND localized density fluctuation in the bulk system that overcomes a threshold ( $N_c$ )



"SuperStretch"  
PICS collaboration  
Lyon-Madrid  
in progress



Muchas gracias!